

Infinite-dimensional stochastic differential equations arising from random matrices

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- Random matrices / Airy RPFs
- Dynamical soft edge scaling / Infinite-dim SDEs of Airy RPFs
- General theorems of ISDEs on strong solutions:
- Dynamical rigidity of 2D Coulomb interacting Brownian motions

- The Gaussian unitary ensembles are Hermitian random matrices

$$M^N = \begin{pmatrix} m_{11} & \cdots & m_{1N} \\ \cdots & \cdots & \cdots \\ m_{N1} & \cdots & m_{NN} \end{pmatrix} \quad m_{ij} = \bar{m}_{ji}.$$

Here $m_{ij} = m_{ij}^1 + \sqrt{-1}m_{ij}^2$. $\{m_{ij}^1, m_{ij}^2, m_{kk}\}_{1 \leq i \leq j \leq N, 1 \leq k \leq N}$ are independent real Gaussian r. v. with mean free, variance 1.

- The distribution of eigen values of M^N is given by ($\beta = 2$)

$$m_\beta^N(d\mathbf{x}_N) = \frac{1}{Z} \prod_{i < j}^N |x_i - x_j|^\beta e^{-\frac{\beta}{4} \sum_{i=1}^N |x_i|^2} d\mathbf{x}_N, \quad (1)$$

- The dist of $N^{-1} \sum_{i=1}^N \delta_{x_i/\sqrt{N}}$ under m_β^N converges to

$$\varsigma(x)dx = \frac{1}{2\pi} \sqrt{4 - x^2} dx \quad (\text{the semi-circle law}) . \quad (2)$$

Airy RPF: $\mu_{\text{Ai},\beta}$

Take the scaling $x_i \mapsto 2\sqrt{N} + s_i N^{-1/6}$ in $m_\beta^N(dx_N)$. Then we have

$$\tilde{\mu}_{\text{Ai},\beta}^N(ds_N) = \frac{1}{Z} \prod_{i<j}^N |s_i - s_j|^\beta e^{-\frac{\beta}{4} \sum_{k=1}^N |2\sqrt{N} + N^{-1/6} s_k|^2} ds_N.$$

Let $\mu_{\text{Ai},\beta}^N = \tilde{\mu}_{\text{Ai},\beta}^N \circ u^{-1}$. ($u(s) = \sum_i \delta_{s_i}$ is the unlabeled map)

Then $\mu_{\text{Ai},\beta}$ is the TDL of $\mu_{\text{Ai},\beta}^N$:

$$\lim_{N \rightarrow \infty} \mu_{\text{Ai},\beta}^N = \mu_{\text{Ai},\beta} \quad (\text{soft-edge scaling})$$

We consider the natural SDE associated with

$$\tilde{\mu}_{\text{Ai},\beta}^N(ds_N) = \frac{1}{Z} \prod_{i<j}^N |s_i - s_j|^\beta e^{-\frac{\beta}{4} \sum_{i=1}^N |2\sqrt{N} + N^{-1/6}s_i|^2} ds_N. \quad (3)$$

We consider the energy form associated with (3). Then

$$\begin{aligned} \int_{\mathbb{R}^N} \frac{1}{2} \sum_{i=1}^N \frac{\partial f}{\partial s_i} \frac{\partial g}{\partial s_i} d\tilde{\mu}_{\text{Ai},\beta}^N &= \\ \int_{\mathbb{R}^N} \frac{1}{2} \left\{ -\Delta f - \sum_{i=1}^N \frac{\partial f}{\partial s_i} \left(\sum_{j \neq i}^N \frac{\beta}{s_i - s_j} - \beta \left\{ N^{\frac{1}{3}} + \frac{s_i}{2N^{\frac{1}{3}}} \right\} \right) \right\} g d\tilde{\mu}_{\text{Ai},\beta}^N \end{aligned}$$

Hence we obtain the SDE of the natural N particle dynamics:

$$dX_t^{N,i} = dB_t^i + \frac{\beta}{2} \left(\left\{ \sum_{j=1, j \neq i}^N \frac{1}{X_t^{N,i} - X_t^{N,j}} \right\} - \left\{ N^{\frac{1}{3}} + \frac{1}{2N^{\frac{1}{3}}} X_t^{N,i} \right\} \right) dt \quad (4)$$

We will detect the limit ISDE of (4) and solve it.

Let $\varrho^N(x) = N^{\frac{1}{3}}\varsigma(xN^{-\frac{2}{3}} + 2)$ be the rescaled semicircle. Then

$$\lim_{N \rightarrow \infty} \varrho^N(x) = \frac{\sqrt{-x}}{\pi} \mathbf{1}_{(-\infty, 0]}(x) =: \varrho(x)$$

$$N^{\frac{1}{3}} = \int_{\mathbb{R}} \frac{\varrho^N(x)}{-x} dx.$$

Hence we have, as $N \rightarrow \infty$,

$$dX_t^{N,i} = dB_t^i + \frac{\beta}{2} \left\{ \left(\sum_{j=1, j \neq i}^N \frac{1}{X_t^{N,i} - X_t^{N,j}} \right) - \left(N^{\frac{1}{3}} + \frac{1}{2N^{\frac{1}{3}}} X_t^{N,i} \right) \right\} dt$$

$$\sim dB_t^i + \frac{\beta}{2} \left\{ \left(\sum_{j=1, j \neq i}^N \frac{1}{X_t^{N,i} - X_t^{N,j}} \right) - \left(\int_{\mathbb{R}} \frac{\varrho^N(x)}{-x} dx \right) \right\} dt$$

$$dX_t^i = dB_t^i + \frac{\beta}{2} \lim_{r \rightarrow \infty} \left\{ \left(\sum_{|X_t^j| < r, j \neq i} \frac{1}{X_t^i - X_t^j} \right) - \left(\int_{|x| \leq r} \frac{\varrho(x)}{-x} dx \right) \right\} dt$$

Thm 1 (O.-Tanemura). Let $\beta = 1, 2, 4$.

Consider the ISDE of $\mathbf{X}_t = (X_t^i)_{i \in \mathbb{N}} \in \mathbb{R}^\infty$:

$$dX_t^i = dB_t^i + \frac{\beta}{2} \lim_{r \rightarrow \infty} \left\{ \left(\sum_{\substack{|X_t^j| < r, \\ j \neq i}} \frac{1}{X_t^i - X_t^j} \right) - \int_{|x| < r} \frac{\varrho(x)}{-x} dx \right\} dt$$

$$\mathbf{X}_0 = \mathbf{s}.$$

(1) The ISDE has a unique, strong solution for $\mu_{\Delta_i, \beta} \circ \ell^{-1}$ -a.s. \mathbf{s} .

Here ℓ is a label, i.e. $\ell : \sum_i \delta_{s_i} \mapsto \mathbf{s} = (s_i) \in \mathbb{R}^\infty$.

(2) Let $X_t = \sum_{i=1}^{\infty} \delta_{X_t^i}$ be the associated unlabeled dynamics.

Then X_t is a $\mu_{\Delta_i, \beta}$ -reversible diffusion.

- $\mathbf{X}_t^N = (X_t^{N,i})_{i=1}^N \in \mathbb{R}^N$ (labeled N particle dynamics):

$$dX_t^{N,i} = dB_t^i + \frac{\beta}{2} \sum_{j=1, j \neq i}^N \frac{1}{X_t^{N,i} - X_t^{N,j}} dt - \frac{\beta}{2} \left\{ N^{\frac{1}{3}} + \frac{1}{2N^{\frac{1}{3}}} X_t^{N,i} \right\} dt$$

- $\mathbf{X}_t = (X_t^i)_{i=1}^\infty \in \mathbb{R}^\infty$ (labeled ∞ particle dynamics):

$$dX_t^i = dB_t^i + \frac{\beta}{2} \lim_{r \rightarrow \infty} \left\{ \left(\sum_{|X_t^j| < r, j \neq i} \frac{1}{X_t^i - X_t^j} \right) - \int_{|x| < r} \frac{\varrho(x)}{-x} dx \right\} dt$$

Thm 2 (O.-Tanemura). Let $\beta = 1, 2, 4$ and $m \in \mathbb{N}$.

Take a label of \mathbf{X}^N and \mathbf{X} in decreasing order: $X_t^i > X_t^{i+1}$ for all i .

Denote by $\mathbf{X}^{N,m}$ and \mathbf{X}^m the first m -components of \mathbf{X}^N and \mathbf{X} .

Suppose $u(\mathbf{X}_0^N) \sim \mu_{\text{Ai},\beta}^N$ and $u(\mathbf{X}_0) \sim \mu_{\text{Ai},\beta}$. Then

$\mathbf{X}^{N,m}$ converge to \mathbf{X}^m weakly in $C([0, \infty); \mathbb{R}^m)$.

- Let $\beta = 1, 2, 4$. Suppose $X_0 \sim \mu_{\text{Ai},\beta}$.

Then the distribution of the top particle X^1 satisfy

$$X_t^1 \sim F_\beta$$

Here F_β is the β Tracy-Widom distribution.

- So far the stochastic dynamics related to Airy RPF was constructed only for $\beta = 2$ by Johansson, Spohn, and others by the method of space-time correlation functions. This stochastic dynamics is same as that of Thm 1.
- Our result is the first time to construct infinite dimensional dynamics for β other than 2.
- When $\beta = 2$, the top particle X^1 is called the Airy process. The Airy process is a scaling limit of various models.

General theorems of ISDEs on strong solutions:

(1) Quasi-Gibbs, (2) Log derivative.

- For an interaction $\Psi : \mathbb{R}^d \times \mathbb{R}^d \rightarrow \mathbb{R} \cup \{\infty\}$

$$dX_t^i = dB_t^i - \frac{\beta}{2} \sum_{j=1, j \neq i}^{\infty} \nabla \Psi(X_t^i, X_t^j) dt \quad (i \in \mathbb{N})$$

- More generally, for $b : \mathbb{R}^d \times S \rightarrow \mathbb{R}^d \cup \{\infty\}$

$$dX_t^i = dB_t^i + b(X_t^i, X_t^{i*}) dt \quad (i \in \mathbb{N})$$

Here X_t^{i*} denotes the unlabeled particles other than X_t^i :

$$X_t^{i*} = \sum_{j \neq i} \delta_{X_t^j}$$

- The ISDEs describe the motion of infinitely many particles consisting of **only one species**. In fact, $b(x, s)$ is independent of i .

- $S_r = \{s \in \mathbb{R}^d; |s| \leq r\}$,
- $S = \{s = \sum_i \delta_{s_i}, s(S_r) < \infty (\forall r)\}$ (configuration space)

- S is the space of **unlabeled** particles (small ∞ dim)
 $(\mathbb{R}^d)^{\mathbb{N}}$ is the space of **labeled** particles (large ∞ dim)

- $\pi_r, \pi_r^c: S \rightarrow S, \pi_r(s) = s(\cdot \cap S_r), \pi_r^c(s) = s(\cdot \cap S_r^c)$.
- For a RPF μ we set the regular conditional probability:

$$\mu_{r,\xi}^m(\cdot) = \mu(\pi_r \in \cdot | s(S_r) = m, \pi_r^c(s) = \pi_r^c(\xi))$$

- Hamiltonian on S_r :

$$\mathcal{H}_r = \sum_{s_i, s_j \in S_r, i < j} \Psi(s_i - s_j)$$

Def: μ is Ψ -quasi-Gibbs measure if $\exists c_{r,\xi}^m$ s.t.

$$c_{r,\xi}^m e^{-\mathcal{H}_r} d\Lambda_r^m \leq \mu_{r,\xi}^m \leq c_{r,\xi}^m e^{-\mathcal{H}_r} d\Lambda_r^m$$

Here Λ_r^m is the Poisson RPF with intensity $1_{S_r} ds$ on $\{s(S_r) = m\}$.

That is, $\Lambda_r^m = \prod_{k=1}^m 1_{S_r}(s_k) ds_k / \text{permutation}$

- Gibbs measures \Rightarrow Quasi-Gibbs measures.
- RPFs with log potentials are not Gibbs measures.
- μ is (canonical) Gibbs measures if

$$\mu_{r,\xi}^m = \frac{1}{\mathcal{Z}} e^{-\mathcal{H}_r - \sum_{s_i \in S_r, \xi_j \notin S_r} \Psi(s_i, \xi_j)} d\Lambda_r^m \quad (5)$$

If $\Psi(x, y) = -\log|x - y|$, then (5) does not make sense.

$$c_{r,\xi}^m e^{-\mathcal{H}_r} d\Lambda_r^m \leq \mu_{r,\xi}^m \leq c_{r,\xi}^m e^{-\mathcal{H}_r} d\Lambda_r^m \quad (\text{quasi-Gibbs property})$$

Lem 1 (O. 13 AOP, O. 13 SPA, O.-Tanemura).

- (1) *Ginibre RPF is a $2 \log |x|$ -quasi Gibbs measure.*
- (2) *Airy $_{\beta}$ RPF are $\beta \log |x|$ -quasi Gibbs m for $\beta = 1, 2, 4$.*
- (3) *Sine $_{\beta}$ RPF are $\beta \log |x|$ -quasi Gibbs m for $\beta = 1, 2, 4$.*

• **Conjecture:** The following is a quasi Gibbs measure

- (1) Airy $_{\beta}$, Sine $_{\beta}$, Bessel $_{\beta}$ for all $0 < \beta < \infty$.
- (2) All determinantal RPFs.
- (3) Zero points of GAFs.

(A1) μ is a Ψ -quasi-Gibbs m with upper-semicont Ψ .

(A2) $\sum_{k=1}^{\infty} k\mu(S_r^k) < \infty$, $\sigma_r^k \in L^2(S_r^k, dx)$.

Here $S_r^k = \{s(S_r) = k\}$, σ_r^k is k -density fun on S_r .

Thm 3 (O.96 CMP, O.13 AOP). Assume (A1) and (A2).

Then \exists μ -reversible diffusion $X_t = \sum_i \delta_{X_t^i}$.



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Thm 4 (O.96 CMP, O.13 AOP). Assume (A1) and (A2).

Then $\exists \mu$ -reversible diffusion $X_t = \sum_i \delta_{X_t^i}$.

We next present an SDE representation of

$$X_t = \sum_{i=1}^{\infty} \delta_{X_t^i}.$$

We detect the ISDE that $X_t = (X_t^i)$ satisfy.

The key notion is “log derivative d^μ of μ ”.

- ρ^1 is 1-correlation function of μ if

$$\int_S \sum_i f(s_i) \mu(ds) = \int_{\mathbb{R}^d} f(x) \rho^1(x) dx \quad \text{for all } f.$$

- Let μ_x be the reduced Palm measure of μ conditioned at x

$$\mu_x(\cdot) = \mu(\cdot - \delta_x \mid s(\{x\}) \geq 1)$$

- Let μ^1 be the 1-Campbell measure on $\mathbb{R}^d \times S$:

$$\mu^1(A \times B) = \int_{A \times B} \rho^1(x) \mu_x(ds) dx$$

- $d^\mu \in L^1_{\text{loc}}(\mathbb{R}^d \times S, \mu^1)$ is called the **log derivative** of μ if

$$\int_{\mathbb{R}^d \times S} f d^\mu d\mu^1 = - \int_{\mathbb{R}^d \times S} \nabla_x f d\mu^1 \quad \forall f \in C_0^\infty(\mathbb{R}^d) \otimes \mathcal{D}_0$$

- Very informally $d^\mu = \nabla_x \log \mu^1$.

(A3) $\{X_t^i\}_{i \in \mathbb{N}}$ do not collide each other

$$\Leftrightarrow \text{Cap}^\mu(\{s \in S; \exists x \in \mathbb{R}^d \text{ s.t. } s(\{x\}) \geq 2\}) = 0$$

(A4) $\int_{\mathbb{R}^d} \rho^1(x) e^{-|x|^{2-c}} dx < \infty$ for some $c > 0$.

(A5) The log derivative $d^\mu \in L_{loc}^1(\mu^1)$ exists.

Thm 5. (O.12(PTRF)) (A1)–(A5) $\Rightarrow \exists S_0 \subset S$ such that

$$\mu(S_0) = 1, \tag{6}$$

and that, for $\forall s \in u^{-1}(S_0)$, there exists (\mathbf{X}, \mathbf{B}) , where $\mathbf{X} = (X^i)$ and $\mathbf{B} = (B^i)$, satisfying

$$dX_t^i = dB_t^i + \frac{1}{2} d^\mu(X_t^i, \sum_{j \neq i} \delta_{X_t^j}) dt, \quad (X_0^i)_{i \in \mathbb{N}} = \mathbf{s} \tag{7}$$

Here $u: S^{\mathbb{N}} \rightarrow S$ such that $u((s_i)) = \sum_i \delta_{s_i}$.

Let $\mathcal{T} = \mathcal{T}(S)$ be the tail σ field of S :

$$\mathcal{T} = \bigcap_{r=1}^{\infty} \sigma[\pi_r^c].$$

Lem 2. \mathcal{T} is μ -trivial for any determinantal RPF μ .

Let $\mu(\cdot|\mathcal{T})$ be the regular conditional probability. Then

$$\mu(\cdot) = \int_S \mu(\cdot|\mathcal{T})(\xi) \mu(d\xi)$$

Lem 3. Let μ be a quasi-Gibbs measures.

Then \mathcal{T} is $\mu(\cdot|\mathcal{T})(\xi)$ -trivial for μ -a.s. ξ .

existence of strong solution

$$dX_t^i = dB_t^i + \frac{1}{2}d^\mu(X_t^i, \sum_{j \neq i} \delta_{X_t^j})dt, \quad (X_0^i)_{i \in \mathbb{N}} = \mathbf{s}$$

Let $d^\mu(x, \mathbf{s}) = \check{d}^\mu(x, \mathbf{s}_m, \mathbf{s}_m^*)$,

where $\mathbf{s}_m = (s_1, \dots, s_m)$, $\mathbf{s}_m^* = \sum_{l > m} \delta_{s_l}$ for $\mathbf{s} = \sum_i \delta_{s_i}$,
and $\check{d}^\mu(x, \mathbf{s}_m, \mathbf{s}_m^*)$ is symmetric in \mathbf{s}_m .

$$H_k^m = \{ \{ (x, \mathbf{s}_m, \mathbf{s}_m^*) \in (\mathbb{R}^d)^{1+m} \times S ;$$

$$|\check{d}^\mu(x, \mathbf{s}_m, \mathbf{s}_m^*) - \check{d}^\mu(\hat{x}, \hat{\mathbf{s}}_m, \mathbf{s}_m^*)| \leq c_k^m |(x, \mathbf{s}_m) - (\hat{x}, \hat{\mathbf{s}}_m)| \}$$

$$H^m = \bigcup_{k=1}^{\infty} H_k^m, \quad H_k^m \subset H_{k+1}^m \dots$$

(A6) $\text{Cap}^{\mu^m}((H^m)^c) = 0$ for all $m = 0, 1, \dots$

i.e. $P_{(x, \mathbf{s}_m, \mathbf{s}_m^*)}((X_t, \mathbf{X}_t^m, \mathbf{X}_t^{m*}) \in H^m \text{ for all } t) = 1$.

Here μ^m is the m -Campbell measure of μ .

Thm 6 (O.-Tanemura). Assume (A1)–(A6). Then:

(1) $\exists S_0$ such that $\mu(S_0) = 1$ and that the ISDE

$$dX_t^i = dB_t^i + \frac{1}{2}d^\mu(X_t^i, \sum_{j \neq i} \delta_{X_t^j})dt,$$

$$(X_0^i)_{i \in \mathbb{N}} = s$$

has **a strong solution** for $s = (s_i) \in S^{\mathbb{N}}$ s.t. $\sum_i \delta_{s_i} \in S_0$.

(2) The ass unlabeled diffusion $X = \sum_i \delta_{X^i}$ satisfies

the absolutely continuity condition:

$$P_{\mu_\xi} \circ X_t^{-1} \prec \mu_\xi \quad (\forall t) \quad \text{for } \mu\text{-a.s. } \xi$$

Here $\mu_\xi = \mu(\cdot | \mathcal{T})(\xi)$ is the tail decomposition in Lem 3.

Uniqueness of strong solutions 1

Thm 7 (O.-Tanemura). Assume (A1)–(A6).

Let $\mathbf{X} = (X^i)$ and $\widehat{\mathbf{X}} = (\widehat{X}^i)$ be strong sol of the ISDE

$$dX_t^i = dB_t^i + \frac{1}{2}d^\mu(X_t^i, \sum_{j \neq i} \delta_{X_t^j})dt, \quad (X_0^i)_{i \in \mathbb{N}} = \mathbf{s} = (s_i)_{i \in \mathbb{N}}$$

on the same Br m. Let $X_t = \sum_i \delta_{X_t^i}$ and $\widehat{X}_t = \sum_i \delta_{\widehat{X}_t^i}$.

Suppose that *the absolutely continuity condition* is satisfied for μ -a.s. ξ :

$$P_{\mu_\xi} \circ X_t^{-1} \prec \mu_\xi \text{ and } P_{\mu_\xi} \circ \widehat{X}_t^{-1} \prec \mu_\xi \quad (\forall t)$$

Then

$$\mathbf{X} = \widehat{\mathbf{X}} \text{ a.s.} \quad \text{for } \mu\text{-a.s. } \mathbf{s} = \sum_{i=1}^{\infty} \delta_{s_i}$$

Cor 1 Assume (A1)–(A6). Assume
(A7) μ is tail trivial.

Then the strong solution $\mathbf{X} = (X^i)$ satisfying

$$P_\mu \circ X_t^{-1} \prec \mu \quad \text{for all } t$$

is unique for μ -a.e. $x = \sum_i \delta_{x_i}$.

- Recall that tail σ -fields of determinantal RPFs, in particular, Airy, Sine, Ginibre RPFs with $\beta = 2$ are trivial.



Cor 1 Assume (A1)–(A6). Assume (A7) μ is tail trivial.

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Cor 2 The ISDE

$$dX_t^i = dB_t^i + b(X_t^i, X_t^{i*})dt \quad (i \in \mathbb{N}) \quad (8)$$

has a strong solution if the differential equation of RPFs μ

$$\frac{1}{2} \nabla_x \log \mu(x, s) = b(x, s) \quad (9)$$

has a solution μ satisfying (A1)–(A4), and (A6).

(The tail σ -field of S is “the boundary condition” of (8), (9)).

Strong solutions of ISDEs: Remarks 1/2

Remark 1 Short history of strong solutions of the ISDEs

$$dX_t^i = dB_t^i - \frac{\beta}{2} \sum_{i=1, j \neq i}^{\infty} \nabla \Psi(X_t^i - X_t^j) dt$$

- Lang ('78, '79, ZWVG) : $\Psi \in C_0^3(\mathbb{R}^d)$, μ is grand canonical Gibbs m. “stationary Markov solutions”. Lippner, Rost (1D)
- Fritz ('87, AOP): $\Psi \in C_0^3(\mathbb{R}^d)$, μ is grand canonical Gibbs m. $d \leq 4$, “Non-equilibrium solutions” and “strong Markov”.
- Tamemura ('96 PTRF): Hard core balls, “strong Markov”
- Rölyly-Tanemura: hard core + “exponential decay Ψ ”
- All authors above used “Ito scheme” in infinite-dimension space. i.e., they use the “Lipschitz continuity” of the coefficients. This makes their proof very hard.

Strong solutions of ISDEs: Remarks 2/2

Remark 2: • Our method is as follows:

- 1) Construct (not strong) solution (\mathbf{X}, \mathbf{B}) by Thm 5.
- 2) Introduce **infinite**-many, **finite**-dimensional SDEs with consistency by using (\mathbf{X}, \mathbf{B}) .
- 3) By (A6), the first m -components \mathbf{X}^m of \mathbf{X} is a unique strong solution of the m 'th finite dimensional SDE.
- 4) \mathbf{X}^m is $\sigma[\mathbf{B}, \mathbf{X}^{m*}]$ -measurable. Here $\mathbf{X}^{m*} = (X^{m+1}, \dots)$
- 5) Hence the limit \mathbf{X} is $\sigma[\mathbf{B}] \vee \bigcap_{m=1}^{\infty} \sigma[\mathbf{X}^{m*}]$ -measurable.
- 6) From the μ -tail triviality of S , we deduce the triviality of the tail σ -field $\bigcap_{m=1}^{\infty} \sigma[\mathbf{X}^{m*}]$ of the labeled path space $C([0, \infty) : (\mathbb{R}^d)^{\mathbb{N}})$.

Examples: Gibbs measures

- All example below satisfy (A1)–(A6).

Gibbs measures :

- All Gibbs measures with Ruelle's class potentials Ψ (smooth outside the origin). Non-collision (A3) does not hold in general. But it always holds for $d \geq 2$ and, for repulsive interaction Ψ in $d = 1$.
- In this case, the SDEs become

$$dX_t^i = dB_t^i - \frac{\beta}{2} \sum_{j \neq i} \nabla \Psi(X_t^i - X_t^j) dt. \quad (10)$$

Lennard-Jones 6-12 potential: Set $d = 3$.

$$\Psi_{6,12}(x) = |x|^{-12} - |x|^{-6}$$

The corresponding ISDE is: ($i \in \mathbb{N}$)

$$dX_t^i = dB_t^i + \frac{\beta}{2} \sum_{j=1, j \neq i}^{\infty} \left\{ \frac{12(X_t^i - X_t^j)}{|X_t^i - X_t^j|^{14}} - \frac{6(X_t^i - X_t^j)}{|X_t^i - X_t^j|^8} \right\} dt .$$

Examples: Sine_β RPF (Dyson's model)–bulk scaling limit

Sine_β RPF: $d = 1, \beta = 1, 2, 4$

$$dX_t^i = dB_t^i + \frac{\beta}{2} \lim_{r \rightarrow \infty} \sum_{|X_t^i - X_t^j| < r, j \neq i} \frac{1}{X_t^i - X_t^j} dt$$

The definition of $\mu = \mu_{\text{sin},\beta}$:

$$\mu_{\text{sin},\beta} = \lim_{N \rightarrow \infty} \mu_{\text{sin},\beta}^N \quad (\text{bulk scaling})$$

where

$$\mu_{\text{sin},\beta}^N = \frac{1}{\mathcal{Z}_N} \left\{ \prod_{i < j}^N |s_i - s_j| \right\}^\beta e^{-\frac{\beta}{4N} \sum_{k=1}^N |s_k|^2} d\Lambda_N.$$

Examples: Ginibre RPF

Ginibre RPF: $d = 2$, $\Psi(x) = -2 \log |x|$.

$$dX_t^i = dB_t^i + \lim_{r \rightarrow \infty} \sum_{|X_t^i - X_t^j| < r, j \neq i} \frac{X_t^i - X_t^j}{|X_t^i - X_t^j|^2} dt. \quad (11)$$

- $\mu_{\text{gin},2}$ is the TD limit of

$$\mu_{\text{gin},2}(ds) = \lim_{N \rightarrow \infty} \frac{1}{\mathcal{Z}_N} \prod_{i=1, i < j}^N |s_i - s_j|^2 \prod_{k=1}^N \frac{1}{\pi} e^{-|s_k|^2} \Lambda_N$$

- $\mu_{\text{gin},2}$ is rotation and translation invariant.



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$$dX_t^i = dB_t^i + \lim_{r \rightarrow \infty} \sum_{|X_t^i - X_t^j| < r, j \neq i} \frac{X_t^i - X_t^j}{|X_t^i - X_t^j|^2} dt. \quad (11)$$

- $\mu_{\text{gin},2}$ is the TD limit of

$$\mu_{\text{gin},2}(ds) = \lim_{N \rightarrow \infty} \frac{1}{\mathcal{Z}_N} \prod_{i=1, i < j}^N |s_i - s_j|^2 \prod_{k=1}^N \frac{1}{\pi} e^{-|s_k|^2} \Lambda_N$$

- $\mu_{\text{gin},2}$ is rotation and translation invariant.

Dynamical rigidity of the Ginibre RPF

Ginibre RPF: $d = 2$, $\Psi(x) = -2 \log |x|$.

$$dX_t^i = dB_t^i + \lim_{r \rightarrow \infty} \sum_{|X_t^i - X_t^j| < r, j \neq i} \frac{X_t^i - X_t^j}{|X_t^i - X_t^j|^2} dt. \quad (11)$$

Thm 8 (O.12 PTRF). *The solution (X^i) of (11) satisfies*

$$dX_t^i = dB_t^i + \lim_{r \rightarrow \infty} \sum_{|X_t^j| < r, j \neq i} \frac{X_t^i - X_t^j}{|X_t^i - X_t^j|^2} dt - X_t^i dt. \quad (12)$$

Remark: • Two ISDEs (11) and (12) have common unique, strong solutions.

• Solutions are trapped very thin subset of $(\mathbb{R}^2)^{\mathbb{N}}$.

Dynamical rigidity of Ginibre RPF 2

Ginibre RPF: $d = 2$, $\Psi(x) = -2 \log |x|$.

$$dX_t^i = dB_t^i + \lim_{r \rightarrow \infty} \sum_{\substack{|X_t^i - X_t^j| < r, \\ j \neq i}} \frac{X_t^i - X_t^j}{|X_t^i - X_t^j|^2} dt. \quad (11)$$

Thm 9. *Each tagged particles are sub diffusive:*

$$\lim_{\varepsilon \rightarrow 0} \varepsilon X_{t/\varepsilon^2}^i = 0 \quad \text{for all } i \in \mathbb{N} \quad (13)$$

weakly in $C([0, \infty) : \mathbb{R}^2)$ in $\mu_{\text{gin},2}$ -measure.

- This result is strikingly different from that of interacting Brownian motions with Ruelle's class potentials. Tagged particles are always diffusive for Ruelle's class potentials for $d \geq 2$.