

Nelson model

1. Fundamental facts
2. Nelson model with UV cutoff
3. Renormalized Nelson model

$$\mathcal{F} = \bigoplus_{n=0}^{\infty} L^2_{sym}(\mathbb{R}^{3n}) \rightarrow \bigoplus_{n=0}^{\infty} \bar{\Phi}_n \quad \underbrace{n=0}_{\mathbb{C}}$$

$$\sum_{n=0}^{\infty} \|\bar{\Phi}_n\|^2 < \infty \Leftrightarrow \bar{\Phi} = \bigoplus_{n=0}^{\infty} \bar{\Phi}_n \in \mathcal{F}, \quad \Omega = \{1, 0, 0, \dots\} \in \mathcal{F}$$

creation, annihilation

$$(a^\dagger(f)\bar{\Phi})^{(n)} = \frac{1}{\sqrt{n!}} \sum_{j=1}^n f(k_j) \bar{\Phi}^{(n-1)}(k_1 \dots \widehat{k}_j \dots)$$

$$(a(f)\bar{\Phi})^{(n)} = \sqrt{(n+1)!} \int f(k) \bar{\Phi}^{(n+1)}(k_1 \dots k_n k) dk$$

$$a(f): \mathcal{F}^{(n)} \rightarrow \mathcal{F}^{(n-1)}, \quad a^\dagger(g): \mathcal{F}^{(n)} \rightarrow \mathcal{F}^{(n+1)}$$

- $[a(f), a^\dagger(g)] = (\bar{f}, g)$

- $a(f)^* = a^\dagger(\bar{f})$

- $\phi(\hat{f}) = \frac{1}{\sqrt{2}} [a^\dagger(\hat{f}) + a(\hat{f})]$ Field op.

$$\widetilde{\hat{f}}(k) = \hat{f}(-k) \quad f \text{ real} \Rightarrow \widetilde{\hat{f}} = \overline{\hat{f}} \quad "$$

- $\overline{\left\{ \prod_{r=1}^n \phi(f_r) \Omega \right\}} = \mathcal{F}^{(n)}$

2nd quantization

Relativistic Schrödinger op $\sqrt{-\Delta + m^2} \in L^2(\mathbb{R}^{3d})$

N-body R-Schrödinger op is $\sum_{j=1}^N \sqrt{-\Delta_j + m^2} \in L^2(\mathbb{R}^{3N})$

$$\sum_{j=1}^N \mathbb{1} \otimes \sqrt{-\Delta_j + m^2} \otimes \dots \otimes \mathbb{1} \otimes \bigotimes_{j=1}^N L^2(\mathbb{R}^3)$$

$h: L^2 \rightarrow L^2$ $\sqrt{j^2 h}$

$$d\Gamma(h) = \bigoplus_{n=0}^{\infty} \sum_{j=1}^n \mathbb{1} \otimes \dots \otimes h \otimes \dots \otimes \mathbb{1}$$

$$\begin{aligned} d\Gamma(a^\dagger(k_1) \dots a^\dagger(k_n)) \Omega &= \sum a^\dagger(k_1) \dots a^\dagger(k_n) \Omega \\ d\Gamma \Omega &= 0 \end{aligned}$$

Let T be contraction on L^2

$$\begin{aligned} \Gamma(T) &= \bigoplus_{n=0}^{\infty} \bigotimes^n T & \Gamma(e^{it} h) & \text{one-parameter} \\ \Gamma(T)\Gamma(S) &= \Gamma(TS) & \therefore \Gamma(e^{it} h) &= e^{it} S \end{aligned}$$

We can see that $S = \left(\sum_{j=1}^n \sqrt{-\Delta_j + m^2} \right)$

$\omega = \sqrt{|k|^2 + m^2}$ multiplication

$$d\Gamma(\omega) = H_f \text{ free field } H$$

$d\Gamma(h) \stackrel{P_f}{=} \text{momentum op.}$

$d\Gamma(\mathbb{1}) = N$ number op

We see that

$$\begin{aligned} [d\Gamma(h), a(t)] &= -a(kt) \\ [d\Gamma(h), a^\dagger(t)] &= a^\dagger(kt) \end{aligned} \quad \ominus$$

- $\|a^\#(f) \bar{\Phi}\|_f \leq \| (H_f + \Delta)^{1/2} \bar{\Phi} \| \|f/w\|$

In particular $H_f + \phi(f)$ is s.a. $> -\infty$
if $f/w \in L^2$.

by Kato-Rellich

- $e^{a^\#(f)}$ unbounded op eg. e^{Ω} coherent vectors

hint

- $e^{a^\#(f)} e^{-tH_f}$ is bdd when $\|f/w\| < \infty$

- BCH-formula

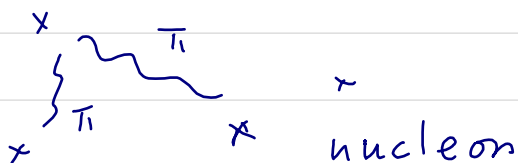
$$e^\phi = e^{\|h\|^2/4} e^{a^\#(-\frac{1}{\sqrt{2}}g)} e^{a^\#(\frac{1}{\sqrt{2}}g)}$$

- $e^{-2t(H_f + \phi(\hat{f}))} = e^{-\frac{1}{4} \int_{-t}^t \int_{-t}^t w(t-s) dt ds} \underbrace{e^{a^\#(U)}}_{e^{-2tH_f}} \underbrace{e^{a^\#(\tilde{U})}}$

where $U = \int_{-t}^t \frac{|f|^2}{w} e^{-|s-t|} w ds$ $\tilde{U} = \int_{-t}^t \frac{|f|^2}{w} e^{-|s-t|} w ds$

$$w(t-s) = \int \frac{|\hat{f}|^2}{w} e^{-|t-s|} w dz$$

Nelson model = Linear interaction between nucleons & pions



$$\left(\sum_{j=1}^N -\frac{1}{2m_j} \sigma_j \right) \otimes \mathbb{1} + \mathbb{1} \otimes H_f + g \sum_{j=1}^N \phi(x_j)$$

$$\phi(x) = \frac{1}{\sqrt{2}} \left\{ a^\dagger \left(\frac{\hat{\psi}}{\sqrt{v}} e^{-ikx} \right) + a \left(\frac{\hat{\psi}}{\sqrt{v}} e^{+ikx} \right) \right\}$$

$$N=1, m_j=1, g=1$$

$$\left(-\frac{1}{2} \sigma + v \right) \otimes \mathbb{1} + \mathbb{1} \otimes H_f + \phi(x)$$

$\hat{\psi}$: UV cutoff function.

$$P_m := -i \nabla_{x_m} \otimes \mathbb{1} + \mathbb{1} \otimes d\Gamma(k_m) \quad \text{Total momentum}$$

Let $v=0$. Then we can see that

$$[P_m, H] = 0 \quad \mathcal{G}(P_m) = \mathbb{R}$$

$$H = \int_{\mathbb{R}^3} H(p) dP, \quad H(p) = \frac{1}{2} (p - p_f)^2 + \phi(1,1) + H_f$$

m J.

F K F

$(B_t)_{t \geq 0}$ } D-BM i.e.

$(\mathcal{X}, \mathcal{B}, \mathbb{W}^x)$ prob space $\mathbb{E}^x[B_t] = 0$

$\mathbb{E}^x[B_t B_s] = t \wedge s$.

$$\mathbb{E}^x[f(B_t)] = \int f(y) p_t(y) dy = \left(e^{-\frac{1}{2} \langle \cdot, \cdot \rangle} f \right) (\mathbb{W}_t)$$

$$\therefore \left(g, e^{-\frac{1}{2} \langle \cdot, \cdot \rangle} f \right) = \int g(\mathbb{W}_t) \mathbb{E}^x[f(B_t)]$$

$$\bullet \left(g, e^{-\frac{1}{2} \langle \cdot, \cdot \rangle + \langle \cdot, v \rangle} f \right) = \int \mathbb{E}^x \left[g(\mathbb{W}_0) + (B_t) \right] \mathbb{E}^x \left[g(\mathbb{W}_t) + (B_t) e^{-\int_0^t \langle v, B_s \rangle ds} \right]$$

$$\left(\Phi, e^{-\frac{1}{2} \langle \cdot, \cdot \rangle + \langle \cdot, v \rangle} \Psi \right) = \frac{1}{2} \int_0^t \left(\Phi, e^{-\frac{1}{2} \langle \cdot, \cdot \rangle + \langle \cdot, v \rangle} \Psi \right) ds$$

$$\bullet \left(g, e^{-\frac{1}{2} \langle \cdot, \cdot \rangle} f \right) = \int \mathbb{E}^x \left[g(\mathbb{W}_{-T}) f(B_T) e^{-\int_{-T}^T \langle v, B_s \rangle ds} \right]$$

$$(B_t)_{t \in \mathbb{R}} \therefore B_{-t} \stackrel{\cdot}{=} B_t, \quad B_t \perp B_s \\ \text{for } t < 0 < s$$

$$\mathbb{E}[f(B_t)] = \int f(y) P_t(x, y) dy$$

$$= \left(e^{-\frac{t}{2}(-\sigma)} f \right)(x) \quad \{z_1, z_2\}$$

$$(f, e^{t h_0} g) = \int dx f(x) \mathbb{E}^x [g(B_t)]$$

$$(f, e^{t(h_0 + v)} g) = \int dx \mathbb{E}^x \left[f(B_0) g(B_t) e^{\int_0^t v(B_s) ds} \right]$$

$$h_0 + v = -\frac{1}{2} \sigma^2 v$$

We can extend this

$$\frac{1}{2} [\sigma \cdot (-i\nu - a)]^2 + v \quad \omega$$

$$\sqrt{[\sigma \cdot (-i\nu - a)]^2 + m^2} + v \quad \dots \text{etc}$$

$$\psi\left(\frac{1}{2} \sigma \cdot (-i\nu - a)\right) + v \quad \psi: \text{Bernstein function}$$

$$\psi(x) = \sqrt{\frac{1}{2} x^2 + m^2}$$

$$\text{Thm} \quad (F, e^{-2\Gamma H} G) = \int d^4x \bar{E}^x \left[(F|_{\beta_T}, \underbrace{e^{-\int_{-t}^t v(\beta_s) ds} e^{i\tilde{U}}}_{G^+(U)}, \underbrace{e^{-2\Gamma H} a(\tilde{U})}_{e^{-2\Gamma H} a(\tilde{U})}) \right]$$

$$S = \int_{-T}^T dt \int_{-T}^T ds W(\beta_s - \beta_T, s-t) \quad \dots \quad W(x+t) = \int \frac{\hat{p}}{\sqrt{2\pi}} e^{-H|w|} e^{-i\kappa\beta_s} ds$$

$$F = f \otimes \Omega, \quad G = g \otimes \Omega \otimes \epsilon$$

$$(F, e^{-2\Gamma H} G) = \int d^4x \bar{E}^x \left[e^{\frac{1}{2}S} f|_{\beta_T} g|_{\beta_T} \right]$$

$$d = 3, \quad h = -\frac{1}{2} \Delta + V$$

$$H = h \otimes \mathbb{1} + \mathbb{1} \otimes H_f + \phi(x)$$

$$\phi(x) = \frac{1}{\sqrt{2}} \left\{ a^+ \left(\frac{\hat{\psi}}{\sqrt{\omega}} e^{-ikx} \right) + a \left(\frac{\tilde{\psi}}{\sqrt{\omega}} e^{+ikx} \right) \right\}$$

$$\tilde{\psi} = \hat{\psi}(-k) \quad \text{if } \psi \text{ is } \mathbb{R}\text{-valued, } \tilde{\psi} = \overline{\hat{\psi}}$$

Thus $\phi(x)^* = \phi(x)$ for each $x \in \mathbb{R}^3$

$$\boxed{\hat{\psi}/\sqrt{\omega} \in L^2}$$

Thm $\hat{\psi}/\omega \in L^2$. Then H is self-adjoint
in $D(h \otimes \mathbb{1}) \cap D(\mathbb{1} \otimes H_f)$

$$\textcircled{!} \quad \| \phi(x) F \| \leq \varepsilon \| (h \otimes \mathbb{1} + \mathbb{1} \otimes H_f) F \| (\| \hat{\psi}/\sqrt{\omega} \| + \| \psi/\omega \|) + b_\varepsilon \| F \| \quad \lrcorner$$

Ground state

$$\begin{cases} \int |\hat{\varphi}|^2 / \omega^3 dx < \infty & \text{IR Infrared regular} \\ \int |\hat{\varphi}|^2 / \omega^2 dx = \infty & \text{IR singular} \end{cases}$$

$$\begin{cases} m > 0 & \text{massive} \\ m = 0 & \text{massless} \end{cases}$$

$$\hat{\varphi}(0) \neq 0 \Rightarrow \begin{cases} \text{massive} \rightarrow \text{IR regular} \\ \text{massless} \rightarrow \text{IR singular} \end{cases}$$

$$\textcircled{1} \quad V(x) = -\frac{1}{|x|}, -\frac{e^{-|x|}}{|x|}, |x|^2 \text{ etc}$$

Assumption V | ① $\lim_{|x| \rightarrow \infty} V(x) = \infty$

② $V(x) \leq -\frac{a}{|x|^\alpha} \quad (0 < \alpha < 2)$

Thm . Suppose that h has a ground state and IR regular $\Rightarrow H$ has a ground state

• IR singular $\Rightarrow H$ has no ground state.

I mentioned that way to make it easier for you all to understand

$m = 0$



1960 - 1970

$m = 0$

IR regular

↙ \exists ground state



1995 ~

- BFS, AH 191<<1
- H. Spohn \forall g 1998
- C. Gérard \forall g 2000
- GLL 2001

IR singular

↙ no ground state



↙ Arai HH 1999

LMS 2001

- h has no ground state, but $1g1 \gg 1$
 H has a ground state Enhanced binding

H + Spohn 2001

- Nelson model on mfd

Gérard + F.H. + Panatier + Suzuki 2011, 12, 13

Nelson on manifolds

$$H_{\text{mfid}} = \frac{1}{2} \partial^\mu g_{\mu\nu}(x) \partial^\nu + V + dP(\hat{\omega}) + \hat{\phi}(x)$$

$$\omega = \sqrt{|k|^2 + m^2} \rightarrow \sqrt{-\partial + m^2}$$

=

$$\rightarrow \sqrt{\partial^\mu g_{\mu\nu} \partial^\nu + m^2(x)}$$

$m^2(x)$



$\hat{\phi}$ variable mass

$$\underline{m^2(x) \sim \langle x \rangle^{-1}}$$



massive



massless

$m > 0 \quad \exists \varphi_m \text{ ground state} \quad \|\varphi_m\| = 1$

$\therefore \lim_m \varphi_m = \varphi_0 \quad \varphi_0 \neq 0$

$\exists \sum \varphi_m \rightarrow 0$

$\varphi_m = \sum_{n=0}^{\infty} \varphi_m^{(n)}$

$\varphi_m^{(n)} \in W^{1,p}(U)$

$i: W^{1,p}(U) \rightarrow L^q$

$U \subset \mathbb{R}^3$

$q < \frac{dP}{d-P}$

$q < \frac{2P}{2-P}$

$\{\varphi_m^{(n)}\}_m \subset L^q$

$[P > 1 \rightarrow q = 2]$

Functional integrators

0) QFT-method Lattice app.

1) Functional integrator

2) $\exists T \text{ cpt op st } T\varphi_m \rightarrow 0 \Rightarrow \varphi_0 \neq 0$

3) $\varphi_m^{(n)} \in W^{1,p}(U) \xrightarrow{i} L^q(U)$

$U \subset \mathbb{R}^{3n}$

cpt

$q < \frac{3P}{3-P}$

Rellich-K.

$P > \frac{6}{5} \rightarrow \{\varphi_m^{(n)}\} \subset L^2$

x Ground state by FH from Springer

ground state expectation (Localization)

$$\varphi = \varphi(x, \phi, n)$$

$$\varphi = \bigoplus_{n=0}^{\infty} \varphi^{(n)} \quad \text{s.t.} \quad \sum_{n=0}^{\infty} \|\varphi^{(n)}\|^2 < \infty$$

$$L^2 \otimes \mathcal{F} = \mathcal{H}$$

$$\sum_{n=0}^{\infty} \|F_n(x)\|_{L^2 \otimes \mathcal{F}^{(n)}}^2$$

$$\|F\|_{\mathcal{H}}^2 = \int \|F(x)\|_{\mathcal{F}}^2 dx = \int \sum_{n=0}^{\infty} \|F_n(x)\|_{\mathcal{F}^{(n)}}^2 dx$$

- $\|\varphi(x)\|_{\mathcal{F}} \sim e^{-|x|}$? $C_1 e^{-C_2|x|} \leq \|\varphi(x)\|_{\mathcal{F}} \leq C_3 e^{-C_4|x|}$
lower bound
- $\sum_{n=0}^{\infty} e^{+\beta n} \|\varphi_n(x)\|_{L^2 \otimes \mathcal{F}}^2 < \infty$? $\forall \beta$
- $\|e^{\beta \phi(g)} \varphi\|_{\mathcal{H}} < \infty \iff \exists \beta > 0 \iff g$

.....

$$\frac{1}{\sqrt{2}} q + i p = a, \quad \frac{1}{\sqrt{2}} q - i p = a^\dagger \quad [a, a^\dagger] = 1$$

$$h = a^\dagger a = \frac{1}{2}(p^2 + q^2 - 1) \quad \text{harmonic oscillator}$$

what is the ground state.

$$\varphi_0 = \pi^{-1/4} e^{-\frac{1}{2}|x|^2}$$

$$\varphi_1 = \alpha \varphi_1 \quad \dots \quad \varphi_n = h_n(x) \varphi_0$$

$$\frac{e^{-t h} f}{\|e^{-t h} f\|} \xrightarrow{f \geq 0} \varphi_0 \quad (t \rightarrow \infty) \quad \because (t, \varphi_0) \neq 0$$

Actually $f = 1$

$$\begin{aligned} e^{-t h} f &= \mathbb{E}^x \left[e^{-\frac{1}{2} \int_0^t |B_s|^2 ds} f(B_t) \right] = \mathbb{E}^x \left[e^{-\frac{1}{2} \int_0^t |B_s|^2 ds} \right] \\ &= e^{-\frac{1}{2} |x|^2 \coth t} \end{aligned}$$

$$\begin{aligned} \frac{e^{-t h} 1}{\|e^{-t h} 1\|^2} &= \frac{e^{-\frac{1}{2} |x|^2 \coth t}}{\left(\int \left(e^{-\frac{1}{2} |x|^2 \coth t} \right)^2 dx \right)^{1/2}} \quad \left(\coth t \rightarrow 1 \right. \\ & \quad \left. t \rightarrow \infty \right) \\ &= \pi^{-1/4} e^{-\frac{1}{2} |x|^2} \quad \checkmark \end{aligned}$$

$$\| e^{\beta |x|^2} \varphi_0 \| \rightarrow \infty \quad \text{as } \beta \uparrow \frac{1}{2}$$

$$\text{QM } \left(-\frac{1}{2}\Delta + V\right)\varphi = E\varphi$$

$$\therefore V(x) \sim |x|^{2n} \quad \therefore \varphi(x) \sim e^{-|x|^{n+1}}$$

$$\begin{aligned} \therefore \varphi &= e^{-tE} \cdot e^{+tE} \varphi = e^{+tE} \mathbb{E}^x \left[e^{-\int_0^t V(B_s) ds} \varphi(B_t) \right] \\ &\leq e^{+tE} \mathbb{E} \left[e^{-\int_0^t V(B_s+x) ds} \right] \|\varphi\|_\infty \end{aligned}$$

$$\sigma_R(x) = \inf \left\{ t > 0 \mid B_t + x \in B_R \right\} \quad \text{Carmona a est}$$

or

Agmon metric:

$$\mathbb{E} \left[\mathbb{1}_{t < \sigma_R(x)} e^{-\int_0^t (V-\bar{E})} \right] \sim e^{-t(V(x)-\bar{E})}$$

$$+ \mathbb{E} \left[\mathbb{1}_{t \geq \sigma_R(x)} e^{-\int_0^t (V-\bar{E})} \right] \sim e^{-\frac{|x|^2}{2t}}$$

$$B_R(x) = I_R(x)$$

$$\frac{|x|^2}{2t} = (V(x) - \bar{E})t \quad \therefore t = \sqrt{\frac{|x|^2}{2(V(x) - \bar{E})}}$$

$$\therefore \varphi(x) \sim e^{-|x| \sqrt{V(x) - \bar{E}}}$$

φ ground state of H .

$$\varphi^T = e^{-TH} f \otimes \Omega / \| e^{-TH} f \otimes \Omega \| \rightarrow \varphi \quad (T \rightarrow \infty) \quad \boxed{f \geq 0}$$

$$\| e^{-\beta H} \varphi \|^2 \rightarrow ? \quad \beta \uparrow ?$$

$N = d\Gamma(\mathbb{1})$ number op.

$$N\varphi = N \bigoplus_{n=0}^{\infty} \varphi_n = \bigoplus_{n=0}^{\infty} n \varphi_n$$

$$\therefore e^{-\beta N} \varphi = \bigoplus e^{-\beta n} \varphi_n$$

$$\text{Hence } (\varphi, e^{-\beta N} \varphi) = \sum (\varphi_n, e^{-\beta n} \varphi_n)$$

$$\lim_{T \rightarrow \infty} (\varphi^T, e^{-\beta N} \varphi^T) = \lim_{T \rightarrow \infty} \frac{(e^{-TH} f \otimes \Omega, e^{-\beta N} e^{-TH} f \otimes \Omega)}{(e^{-TH} f \otimes \Omega, e^{-TH} f \otimes \Omega)}$$

$$= \frac{\int_{\mathbb{H}} \left[e^S f(B_T) f(B_T) e^{-\frac{(1-e^{-\beta})}{T} \int_0^T W(s-t, B_T)} \right] dx}{\int_{\mathbb{H}} \left[e^S f(B_T) f(B_T) \right] dx}$$

$$\int_{\mathbb{H}} \left[e^S f(B_T) f(B_T) \right] dx \quad \text{x-indep}$$

$$\Rightarrow \frac{\mathbb{E}^0 \left[\int e^{S(x)} f(B_T+x) f(B_T+x) dx e^{-\frac{(1-e^{-\beta})}{T} \int W} \right]}{\mathbb{E}^0 \left[\int \dots dx \right]}$$

$$\mathbb{E}^0 \left[\int \dots dx \right]$$

$$\int \frac{-(1 - e^{-\beta})}{e} \int_{-T}^0 \int_0^T W \, d\mu_T \quad \mathcal{H}$$

$$\bigcup_{s > 0} \mathcal{G}(B_t \mid t \in [s, S]) = \mathcal{B}$$

Thm (FH 2014)

μ_T is a measure on $(\mathcal{H}, \mathcal{G}(B))$

s.t. $\exists \mu_\infty$ s.t. $\mu_T(A) \rightarrow \mu_\infty(A) \quad \forall A \in \mathcal{G}(B)$

$$\left| \int_{-T}^0 \int_0^T W \right| \leq \int \frac{|\hat{\varphi}|^2}{W^3} < \infty \quad \swarrow \text{we assume}$$

$$\therefore \int_{-\infty}^0 \int_0^\infty W < \infty$$

$$\therefore \lim_T (\varphi^T, e^{-\beta N} \varphi^T) = (\varphi, e^{-\beta N} \varphi)$$

$$\begin{aligned} &= \sum e^{-\beta n} \|\varphi_n\|^2 \\ &\int_{\mathcal{H}} \frac{-(1 - e^{-\beta})}{e} \int_{-\infty}^0 \int_0^\infty W \, d\mu_\infty \quad \beta \in (\Gamma_0, \infty) \mapsto \beta \in \mathbb{C} \end{aligned}$$

$$\therefore \sum e^{+\beta n} \|\varphi_n\|^2 = \int \frac{-(1 - e^{+\beta})}{e} \int_{-\infty}^0 \int_0^\infty W \, d\mu_\infty < \infty$$

$$\begin{aligned}
 (\varphi, e^{-\beta \phi(g)^2} \varphi) &= \lim_T (\varphi^T, e^{-\beta \phi(g)^2} \varphi^T) \\
 &= \frac{1}{\sqrt{1 + \beta \|g\|^2}} \int_{\mathbb{R}^n} \left[e^{\frac{-\beta K^2}{1 + \beta \|g\|^2}} \right]
 \end{aligned}$$

where

$$K = \frac{1}{2} \int_{-\infty}^{\infty} \left(e^{-|s|\omega} e^{-i\omega \cdot B_s} \hat{\varphi} / \sqrt{\omega}, g \right)_{L^2} ds$$

$$\therefore \| e^{+\frac{\beta}{2} \phi(g)^2} \varphi \|^2 \rightarrow \infty \quad \beta \rightarrow \frac{1}{\|g\|^2}$$

$$H_\varepsilon = \left(-\frac{1}{2}\Delta + V\right) \otimes \mathbb{1} + \mathbb{1} \otimes H_f + \phi_\varepsilon(x)$$

$$\phi_\varepsilon(x) = \frac{1}{\sqrt{2}} \left\{ a^\dagger \left(\frac{e^{-\varepsilon|k|^2/2}}{\sqrt{\omega}} e^{-ik \cdot x} \right) + a \left(\frac{e^{-\varepsilon|k|^2/2}}{\sqrt{\omega}} e^{+ik \cdot x} \right) \right\}$$

$$\varepsilon > 0 \quad \omega = \sqrt{|k|^2 + m^2} \quad m \geq 0$$

Thm (E. Nelson 1964)

$$\text{Let } E_\varepsilon = -\frac{1}{2} \int \frac{e^{-\varepsilon|k|^2}}{\omega} \cdot \frac{1}{|k|^{3/2} + \omega} dk \rightarrow -\infty \quad (\varepsilon \downarrow 0)$$

$$H_\varepsilon - E_\varepsilon \rightarrow \exists H_{\text{ren}} \quad (\varepsilon \downarrow 0) \text{ in semi group.}$$

2014 Gubinelli - H-Lövinzi

we had proven it by path measure.

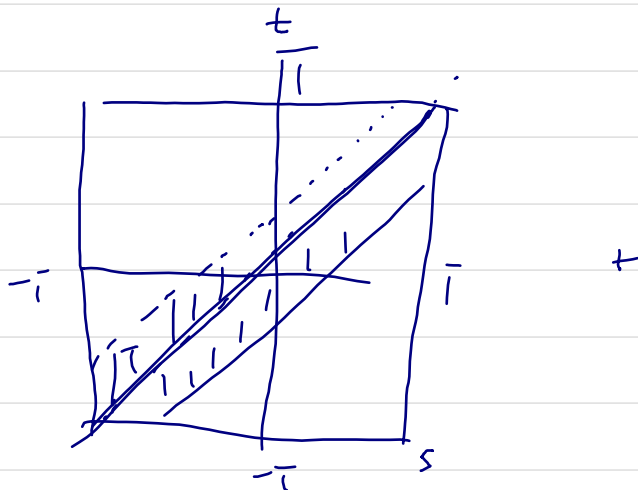
$$V = 0$$

$$(f \otimes \Omega, e^{-2iH_\varepsilon} g \otimes \Omega) = \int dx \int_{\mathbb{T}} \left[e^{\frac{1}{2}S} f(B_{-T}) \int(B_{-T}) \right]$$

where

$$S = \int_{-T}^T \int_{-T}^T \left(\int \frac{e^{-\varepsilon|k|^2}}{2\omega} e^{-|t-s|\omega} e^{-ik(B_t - B_s)} dk \right) dt ds$$

$$\varepsilon \rightarrow 0$$



$$\int_{-T}^T ds \int_s^{s+T} \omega dt = 2T \rho(0,0) - \int_{-T}^T \rho(B_s - B_{-T}, s - T) ds$$

$$+ \int_{-T}^T \int_s^{s+T} \nabla \rho(B_s - B_{-T}, s - t) dB_t ds$$

$$\rho = \rho_\varepsilon(x, t) = \int \frac{e^{-\varepsilon|k|^2} e^{-ikx} e^{-|t|\omega}}{2\omega} dk$$

$$\therefore (f \otimes \Omega, e^{-2T(H+P(00))} g \otimes \Omega)$$

$$= \int dx \mathbb{E} \left[\overline{f(B_T)} g(B_T) e^{\frac{1}{2} S_{\text{nem}}^\varepsilon} \right] \rightarrow \varepsilon \downarrow 0 \text{ conv.}$$

STEP 2 $(f \otimes \Phi, e^{-2T(H+P(00))} g \otimes \Phi)$
 $\rightarrow \text{conv.}$

$$F = F(\phi(t_1) \dots \phi(t_m)) \quad F \in \mathcal{F}(\mathbb{R}^m)$$

STEP 3 $(F, e^{-2T(H+P(00))} G)$
 $\forall F, G \quad ?$

Thm (Key thm) $\frac{H+P(00)}{\varepsilon} > -c \leftarrow \varepsilon\text{-indep}$

$$\forall F, \forall G \quad \exists t_n \quad G_n \in D \text{ dense}$$

$$(F, e^{-2T H} G) \text{ limiting argument}$$

$$\therefore \exists H_{\text{nem}} \text{ st } \lim_n (F, e^{-2T H_{\varepsilon+P(00)}} G)$$

(Matthie Møll 2017)

$$\left(F, e^{-2T H_{ren}} G \right)_{-\int_V} \\ = \int dx \mathbb{E} \left[e^{\frac{1}{2} S_{ren}} \left(F(B_T), e^{a^+(U)} e^{-2T H_f} e^{a(\tilde{U})} G(B_T) \right) \right]$$

S_{ren} = complicated which include stochastic int.

$$U = -\frac{1}{\sqrt{2}} \int_{-T}^T \frac{1}{\sqrt{w}} e^{-|s-T|w} e^{-ik \cdot B_s}$$

$$\tilde{U} = -\frac{1}{\sqrt{2}} \int_{-T}^T \frac{1}{\sqrt{v}} e^{-|s+T|w} e^{+ik \cdot B_s}$$