## CAP FORMS ON $U_{E/F}(4)$

#### TAKUYA KONNO

# 1. What are CAP forms ?

The term CAP in the title is a short hand for the phrase "Cuspidal automorphic representations which Associated to Parabolic subgroups". More precisely, let G be a connected reductive group defined over a number field F, and  $G^*$  be its quasisplit inner form. An irreducible automorphic representation  $\pi = \bigotimes_v \pi_v$  of  $G(\mathbb{A})$  is a *CAP form* if

- (i)  $\pi$  is cuspidal;
- (ii) There exists a residual discrete automorphic representation  $\pi^* = \bigotimes_v \pi_v^*$  of  $G^*(\mathbb{A})$  such that  $\pi_v \simeq \pi_v^*$  at all but finite number of v.

Assuming the suitable generalization of Ramanujan's conjecture, the CAP forms should be exactly the non-tempered cusp forms. Thus such forms should be beautifully described in terms of Arthur's conjecture.

We write  ${}^{L}G = \widehat{G} \rtimes_{\rho_{G}} W_{F}$  for the *L*-group of *G* where  $W_{F}$  is the Weil group of  $\overline{F}/F$ ,  $\overline{F}$  being an algebraic closure of *F*. For the purpose of stating the conjecture, we introduce the hypothetical Langlands group  $\mathcal{L}_{F}$  of *F*. An *A*-parameter for *G* is a continuous homomorphism  $\psi : \mathcal{L}_{F} \times SL(2, \mathbb{C}) \to {}^{L}G$  such that

(1) The following diagram commutes:

$$\mathcal{L}_F \times SL(2, \mathbb{C}) \xrightarrow{\psi} {}^LG$$

$$\downarrow \text{proj} \qquad \qquad \qquad \downarrow \text{proj}$$

$$W_F = W_F$$

- (2)  $\psi|_{\mathcal{L}_F}$  is semisimple and  $\psi(\mathcal{L}_F)$  is bounded;
- (3)  $\psi|_{SL(2,\mathbb{C})}$  is analytic.

We should also have some relevance condition (see Ex. 1.2 below). The set  $\Psi(G)$  of  $\widehat{G}$ conjugacy classes of A-parameters should parameterize the irreducible representations of  $G(\mathbb{A})$  which appear in the  $L^2$ -automorphic spectrum of G.

The A-parameters  $\psi$  associated to CAP forms  $\pi = \bigoplus_{v} \pi_{v}$  can be characterized as follows. The cuspidality of  $\pi$  implies, in particular, it appears in the discrete spectrum, so that  $\psi$  should be *elliptic*. That is,  $\operatorname{Im}\psi$  is not contained in the *L*-group of any proper *F*-parabolic subgroup of *G*. We write  $\Psi_{0}(G)$  for the set of elliptic elements in  $\Psi(G)$ . On the other hand the condition "associated to parabolics" is equivalent to the non-triviality of  $\psi|_{SL(2,\mathbb{C})}$ .  $\Psi_{CAP}(G)$  denotes the subset of  $\psi \in \Psi(G)$  satisfying these two conditions. Let us look at some examples.

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**Example 1.1.** Consider G = GL(n). <sup>*L*</sup>G is the direct product of  $\widehat{G} = GL(n, \mathbb{C})$  and  $W_F$ , and an A-parameter is an n-dimensional representation of  $\mathcal{L}_F \times SL(2, \mathbb{C})$ .  $\psi \in \Psi_0(G)$  is equivalent to its irreducibility. Thus each  $\psi \in \Psi(G)$  is of the form

$$\psi = \varphi \otimes \rho_d,$$

where  $\varphi$  is an m-dimensional irreducible representation of  $\mathcal{L}_F$ ,  $\rho_d$  is the d-dimensional irreducible representation of  $SL(2,\mathbb{C})$  and n = dm.  $\varphi$  should correspond to an irreducible cuspidal representation  $\pi$  of  $GL(m,\mathbb{A})$ . Then it is a result of Moeglin-Waldspurger that the only automorphic representation attached to this is the global Langlands quotient

$$\Pi_{\psi} = \bigotimes_{v} J_{P_{m^{d}}}^{GL(n)}(\pi_{v} | \det |_{v}^{(d-1)/2} \otimes \pi_{v} | \det |_{v}^{(d-3)/2} \otimes \cdots \otimes \pi_{v} | \det |_{v}^{(1-d)/2}),$$

which appears in the residual spectrum. Here  $P_{m^d}$  denotes the standard parabolic subgroup of GL(n) with a Levi component isomorphic to  $GL(m)^d$ . Thus there are no CAP forms on GL(n).

**Example 1.2.** Next look at the multiplicative group  $G = D^{\times}$  of a central division algebra D of dimension  $n^2$  over F.  $\psi \in \Psi_{CAP}(G)$  is of the same form as in the previous example. But there should be a relevance condition on  $\psi$  so that only the parameters of the form

$$\psi = \chi \otimes \rho_n, \quad \dim \chi = 1$$

is relevant.  $\chi : \mathcal{L}_F \to W_F \to \mathbb{C}^{\times}$  is identified with an idele class character of F. Certainly the corresponding representation  $\Pi_{\psi} = \chi \circ \nu_{D/F}$  is a CAP form, where  $\nu_{D/F}$  is the reduced norm of D.

**Example 1.3.** Finally we describe  $\Psi_{CAP}(G)$  for  $G = G_n := U(n)_{E/F}$ , the quasisplit unitary group in n-variables attached to a quadratic extension E of F. We write  $\sigma$  for the generator of Gal(E/F), and define  $\theta_n$  by

$$\theta_n(g) := \operatorname{Ad}\left(\begin{pmatrix} & & 1 \\ & -1 & \\ (-1)^{n-1} & & \end{pmatrix}\right)^t g^{-1}.$$

We realize  $G_n$  so that

$$G_n(R) := \{ g \in GL(n, R \otimes_F E) \, | \, \theta_n({}^{\sigma}g) = g \}$$

for any commutative F-algebra R.  ${}^{L}G_{n} = GL(n, \mathbb{C}) \rtimes_{\rho_{G_{n}}} W_{F}$ , where

$$\rho_{G_n}(w) = \begin{cases} \text{id} & \text{if } w \in W_E; \\ \theta_n & \text{otherwise.} \end{cases}$$

 $\psi$  is known to be determined by its restriction to  $\mathcal{L}_E \times SL(2, \mathbb{C})$ . Now each  $\psi \in \Psi_{CAP}(G)$  is of the form

$$\psi|_{\mathcal{L}_E \times SL(2,\mathbb{C})} = \bigoplus_{i=1}^r (\omega_i \varphi_i) \otimes \rho_{d_i}.$$

Here,

- $\varphi_i$  is an irreducible  $m_i$ -dimensional representation of  $\mathcal{L}_E$ , which corresponds to an irreducible cuspidal representation  $\pi_i$  of  $GL(m_i, \mathbb{A}_E)$  such that  ${}^{\sigma}\pi \simeq \pi^{\vee}$  and  $\omega_{\pi}|_{\mathbb{A}^{\times}}$ , the central character of  $\pi$  restricted to  $\mathbb{A}^{\times}$ , is trivial.
- $\omega_i$  is an idele class character of E whose restriction to  $\mathbb{A}^{\times}$  equals  $\omega_{E/F}^{n-d_i-m_i+1}$ .  $\omega_{E/F}$  is the quadratic character of  $\mathbb{A}^{\times}/F^{\times}$  associated to E/F by the classfield theory.
- $n = \sum_{i=1}^{r} d_i m_i$ , and the  $\omega_i \varphi_i \otimes \rho_{d_i}$  are not equivalent to each other.

What are the CAP forms associated to the A-parameters in Ex. 1.3 ? The cases  $n \leq 2$  is trivial. When n = 3, it was proven by Gelbart-Rogawski that such representations are the theta liftings from  $G_1 = U(1)_{E/F}$ . Today we shall report the result of the joint work with Kazuko Konno in the case n = 4. There are 6 types of parameters for  $G_4$  according to the following list of  $\{(d_i, m_i)\}_i$ :

- (1) Stable parameters:  $(1.a) \{(4,1)\}, (1.b) \{(2,2)\}.$
- (2) Endoscopic parameters: (2.a)  $\{(3,1), (1,1)\}, (2.b) \{(2,1), (1,2)\}, (2.c) \{(2,1), (2,1)\}, (2.d) \{(2,1), (1,1), (1,1)\}.$

In this talk, we shall concentrate to the interesting cases (2.b) and (2.c).

# 2. Local theory

Let us write  $\eta$  and  $\eta'$  for characters of  $\mathbb{A}_E^{\times}/E^{\times}$  whose restriction to  $\mathbb{A}^{\times}$  is trivial. Then the parameters to be considered are given by

$$(2.b) \ \psi_{\eta,\pi}|_{\mathcal{L}_E \times SL(2,\mathbb{C})} = (\eta \otimes \rho_2) \oplus \varphi_{\pi}, \quad (2.c) \ \psi_{\underline{\eta}}|_{\mathcal{L}_E \times SL(2,\mathbb{C})} = (\eta \otimes \rho_2) \oplus (\eta' \otimes \rho_2),$$

where  $\underline{\eta} = (\eta, \eta')$  and  $\pi$  is an irreducible cuspidal representation of  $GL(2, \mathbb{A}_E)$  satisfying the condition of Ex. 1.3. At each place v of F, these give rise to the local parameters  $\psi_v : \mathcal{L}_{F_v} \times SL(2, \mathbb{C}) \to {}^LG_v$ . In this section, we shall construct the local A-packet  $\Pi_{\psi_v}(G)$ attached to such  $\psi_v$ . Again for brevity, we consider only those v where  $E_v := E \otimes_F F_v$ is a quadratic extension of  $F_v$  and  $\pi_v$  is (super) cuspidal in (2.b) and  $\eta_v \neq \eta'_v$  in (2.c). Also to keep in accordance with the theme of this conference, we restrict ourselves to non-archimedean v. (The archimedean case is more explicit and needs some case-by-case treatment.)

2.1. **Parabolic induction.** Recall some assertion from Arthur's conjecture. For  $\psi_v \in \Psi_0(G_v)$  we set  $S_{\psi_v}(G) := \operatorname{Cent}(\psi_v, \widehat{G})/Z(\widehat{G})^{\Gamma_v}$ , where  $\Gamma_v = \operatorname{Gal}(\overline{F_v}/F_v)$ . In the case of  $G = G_n$ , we have  $S_{\psi_v}(G) \simeq (\mathbb{Z}/2\mathbb{Z})^{r-1}$  where r is as in Ex. 1.3. Moreover since F is non-archimedean, we may postulate that there exists a perfect duality

$$\langle , \rangle : \mathcal{S}_{\psi_v}(G) \times \Pi_{\psi_v}(G) \longrightarrow \mathbb{C}^{\times}.$$

Next set

$$\mu_{\psi}: \mathbb{G}_m(\mathbb{C}) \ni t \longmapsto \begin{pmatrix} t \\ t^{-1} \end{pmatrix} \in SL(2, \mathbb{C}) \xrightarrow{\psi} \widehat{G}.$$

There is an *F*-parabolic subgroup  $P_{\psi} = M_{\psi}U_{\psi}$  such that  $\widehat{M}_{\psi} = \text{Cent}(\mu_{\psi}, \widehat{G})$  and  $\mu_{\psi}$  is a  $\widehat{P}_{\psi}$ -dominant cocharacter of  $\widehat{M}_{\psi}$ . Associated to  $\psi_{v}|_{\mathcal{L}_{F_{v}}} \in \Psi(M_{\psi})$  is the tempered *L*-packet  $\prod_{\psi_{v}|\mathcal{L}_{F_{v}}}(M_{\psi})$ . Then Arthur imposed

$$\Pi'_{\psi_v}(G) := \{ J^G_{P_{\psi}}(\tau \otimes e^{\mu_{\psi}}) \, | \, \tau \in \Pi_{\psi_v \mid \mathcal{L}_{F_v}}(M_{\psi}) \} \subset \Pi_{\psi_v}(G).$$

Here, again  $J_P^G(\tau \otimes e^{\lambda})$  denotes the Langlands quotient of  $I_P^G(\tau \otimes e^{\lambda}) = \operatorname{ind}_{P(F_v)}^{G(F_v)}[(\tau \otimes e^{\lambda}) \otimes 1_{U(F_v)}].$ 

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We now examine these constructions in our  $G_4$ -case. The standard parabolic subgroups of  $G_4$  are G,  $P_i = M_i U_i$  (i = 1, 2) and the upper triangular Borel subgroup  $\mathbf{B} = \mathbf{TU}$ , where

$$\begin{split} M_1 &= \left\{ \begin{array}{c|c} m_1(a,g) = \left( \begin{array}{c|c} a & & \\ \hline g & \\ \hline g & \\ \hline \end{array} \right) \left| \begin{array}{c} a \in \operatorname{Res}_{E/F} \mathbb{G}_m \\ g \in G_2 \end{array} \right\}, \\ U_1 &= \left\{ \begin{array}{c|c} u_1(y,\beta) = \left( \begin{array}{c|c} \frac{1}{y''} & y' & \beta - \langle y, y \rangle / 2 \\ \hline 1 & -^{\sigma} y' \\ \hline 1 & 0 \end{array} \right) \left| \begin{array}{c} y = (y'',y') \in W_1 \\ \beta \in \mathbb{G}_a \end{array} \right\}, \\ M_2 &= \left\{ \begin{array}{c|c} m_2(a) = \left( \begin{array}{c|c} a \\ \hline \theta_2(^{\sigma} a) \end{array} \right) \left| \begin{array}{c} a \in H_2 \\ f(^{\sigma} b) = -b \end{array} \right\}, \\ U_2 &= \left\{ \begin{array}{c|c} u_2(b) = \left( \begin{array}{c|c} \frac{1}{2} & b \\ \hline 1 & 1 \end{array} \right) \right| \begin{array}{c} b \in (\operatorname{Res}_{E/F} \mathbb{M}_2) \\ f(^{\sigma} b) = -b \end{array} \right\}. \end{split}$$

Here  $(W_1, \langle , \rangle)$  is the hyperbolic skew-Hermitian space  $(E^2, (-1^1))$  over E.

**Proposition 2.1.** The group  $S_{\psi_v}(G)$  and  $\Pi'_{\psi_v}(G)$  for parameters (2.b), (2.c) are given as follows.

A-parameter	$\mathcal{S}_{\psi_v}(G)$	$\Pi'_{\psi}(G)$
(2.b) $\psi_{\pi_v,\eta_v}$	$\mathbb{Z}/2\mathbb{Z}$	$\{J_{P_1}^G(\eta_v   \mid_{E_v}^{1/2} \otimes \tau_v)   \tau_v \in T_v\}$
(2.c) $\psi_{\underline{\eta}_v}$	$\mathbb{Z}/2\mathbb{Z}$	$\{J^G_{P_2}(I(\eta_v\otimes\eta'_v) \det ^{1/2}_{E_v})\}$

 $T_v$  in (2.b) is the L-packet of  $G_2(F_v)$  whose standard base change to  $GL(2, E_v)$  is  $\pi_v$ .  $I(\eta_v \otimes \eta'_v)$  in (2.c) is the obvious notation for the principal series representation of  $GL(2, E_v)$ .

2.2. Local theta correspondence. We still need to construct the members of  $\Pi_{\psi_v}(G) \setminus \Pi'_{\psi_v}(G)$ . For this we use the theta correspondence. We write  $G_2^+ := G_2$  and  $G_2^-$  for its anisotropic inner form. Fix a non-trivial additive character  $\psi_{F_v}$  of  $F_v$ . For a character  $\eta_v$  of  $E_v^{\times}$  having the trivial restriction to  $F_v^{\times}$ , we can construct the Weil representation  $(\omega_{\eta_v,\pm}, S)$  of the unitary dual pair  $G_2^{\pm}(F_v) \times G_{2n}(F_v)$ . Again for simplicity, we assume that the residual characteristic of  $F_v$  is odd. Then for an irreducible admissible representation  $\pi$  of  $G_{2n}(F_v)$  (resp.  $\tau$  of  $G_2^{\pm}(F_v)$ ), we have its (possibly zero) local Howe correspondent  $\theta_{2n,\eta_v}^{\pm}(\pi)$  (resp.  $\theta_{\eta_v}^{2n}(\tau)$ ), an irreducible representation of  $G_2^{\pm}(F_v)$  (resp.  $G_{2n}(F_v)$ ). Then we can prove the following.

**Proposition 2.2** (local  $\theta$ -correspondence for U(2)). Consider the case n = 1. (i)  $\theta_{2,\eta_v}^{\epsilon}(\pi)$  does not vanish if and only if  $\epsilon(1/2, \pi \times \eta_v, \psi_{F_v})\omega_{\pi}(-1)\lambda(E_v/F_v, \psi_{F_v})^2 = \epsilon$ . (ii) If this is the case, we have

$$\theta_{2,\eta_v}^{\epsilon}(\pi) = \begin{cases} \pi^{\vee} & \text{if } \epsilon = 1, \\ \operatorname{JL}(\pi^{\vee}) & \text{otherwise.} \end{cases}$$

Here  $\pi^{\vee}$  is the contragredient of  $\pi$  and  $JL(\pi^{\vee})$  is the "Jacquet-Langlands correspondent" of  $\pi^{\vee}$ .

**Remark 2.3.** This is the  $\epsilon$ -dichotomy property which is proved by Harris-Kudla-Sweet for general unitary dual pairs. But their result uses the  $\epsilon$ -factor defined by the doubling method of Piatetskii-Shapiro-Rallis. The comparison conjecture between their  $\epsilon$ -factor and that defined by the Langlands-Shahidi method is not yet established in the present case. This is the reason why we rely on the construction of M. Harris using the Shimizu-Jacquet-Langlands correspondence.

In general, we rarely have a good understanding of the ramified local theory of theta correspondences. Except for few examples such as Waldspurger's study on the Shimura correspondence, all such examples put some strong restriction on the representations treated (e.g. trivial in the example of Siegel-Weil formula, and one-dimensional in the study of Shalika-Tanaka).

As was remarked above, we deduce the proposition from the similar result for GU(2) obtained by M. Harris. His method deduces the ramified local theory from the Shimizu-Jacquet-Langlands correspondence combined with the result of Waldspurger-Tunnel-Saito-Prasad on the restriction of a representation of  $GL(2, F_v)$  to elliptic Cartan subgroups. Our task is to deduce the unitary group case from the similitude group setting. This is achieved with the help of the following lemma. Let us identify the unitary similitude group  $\widetilde{G}_2 = GU(2)$  with  $(\operatorname{Res}_{E/F}\mathbb{G}_m \times GL(2))/\mathbb{G}_m$  by

$$(\operatorname{Res}_{E/F}\mathbb{G}_m \times GL(2))/\mathbb{G}_m \ni (z,g) \longmapsto z^{-1}g \in \widetilde{G}_2 \subset \operatorname{Res}_{E/F}GL(2).$$

Thus each irreducible representation  $\widetilde{\pi}$  of  $\widetilde{G}_2(F_v)$  is of the form  $\omega \otimes \pi$ , where  $\omega$  is a character of  $E_v^{\times}$  and  $\pi$  is an irreducible representation of  $GL(2, F_v)$  such that  $(\omega|_{F_v^{\times}})\omega_{\pi} = \mathbf{1}$ .

**Lemma 2.4.** We have  $\widetilde{\pi}|_{G_2(F_v)} = \bigoplus_{\tau \in T} \tau$ , where T is the unique L-packet of  $G_2(F_v)$ whose standard base change lift is  $\omega(\det)\pi_{E_v}$ .  $\pi_{E_v}$  denotes the base change of  $\pi$  to  $GL(2, E_v)$ .

We are now ready to complete the A-packets  $\Pi_{\psi_v}$ . Consider the following diagram of local theta correspondences.

$$\begin{array}{cccc}
G_2^{-\epsilon}(F_v) & \xrightarrow{\theta_{\eta_v}^4} & G_4(F_v) \\
& & & & \\
JL \uparrow & \swarrow & | \\
G_2^{\epsilon}(F_v) & \xleftarrow{\theta_{2,\eta_v}^\epsilon} & G_2(F_v)
\end{array}$$

The left vertical arrow is the "Jacquet-Langlands correspondence", while the right vertical line indicates the Witt tower. For  $\tau_v \in T_v$  in Prop. 2.1, write

$$\epsilon := \epsilon (1/2, \pi \times \eta_v, \psi_{F_v}) \omega_\pi (-1) \lambda (E_v/F_v, \psi_{F_v})^2$$

so that  $\theta_{2,\eta_v}^{\epsilon}(\tau_v) \neq 0$ . The tower property of theta correspondence implies

$$heta_{\eta_v}^4 \circ heta_{2,\eta_v}^\epsilon( au_v) \simeq \pi_{\psi,v}^+ := J_{P_1}^G(\eta_v) \mid_{E_v}^{1/2} \otimes au_v).$$

On the other hand, we know from Prop. 2.2 that  $\theta_{\eta_v}^2 \circ JL \circ \theta_{2,\eta_v}^{\epsilon}(\tau_v)$  is zero. This again combined with the tower property says that

$$\pi_{\psi,v}^{-} := \theta_{\eta_v}^4 \circ \mathrm{JL} \circ \theta_{2,\eta_v}^{\epsilon}(\tau_v)$$

is a cuspidal representation of  $G_4(F_v)$ . We argue similarly in the case (2.c) and obtain the following. **Theorem 2.5.**  $\Pi_{\psi_v}(G) = \{\pi_{\psi,v}^{\pm}\}, \text{ where } \pi_{\psi,v}^{\pm} \text{ are defined above in the case (2.b) and }$ 

$$\pi_{\psi,v}^+ = J_{P_2}^G(I(\eta_v \otimes \eta'_v) |\det|_{E_v}^{1/2}), \quad \pi_{\psi,v}^- := \theta_{\eta_v}^4((\eta_v {\eta'_v}^{-1})_{G_2^-}),$$

in the case (2.c). Here  $(\eta_v {\eta'_v}^{-1})_{G_2^-} : G_2^-(F_v) \xrightarrow{\det} G_1(F_v) \ni x\sigma(x)^{-1} \mapsto (\eta_v {\eta'_v}^{-1})(x) \in \mathbb{C}^{\times}$ .

**Remark 2.6.** (i) Similar results hold for other types of  $\pi_v$  in the case (2.b). In particular, when  $\pi_v$  is the special representation,  $\pi_{\psi,v}^-$  coincides with that in the case (2.c). (ii) The archimedean case is treated in a similar way. In the case (2.b) with  $T_v$  is endoscopic,  $|T_v| = 2$  while its Jacquet-Langlands correspondent is a singleton. But since  $\epsilon(1/2, \tau_v \times \eta_v, \psi_{\mathbb{R}})\omega_{\tau_v}(-1)\lambda(\mathbb{C}/\mathbb{R}, \psi_{\mathbb{R}})^2 = 1$  in this case, the theorem is still valid.

## 3. GLOBAL THEORY

We now define the duality (multiplicity pairing)  $\langle , \rangle : \mathcal{S}_{\psi_v}(G) \times \Pi_{\psi_v}(G) \to \{\pm 1\}$  so that  $\pi_{\psi,v}^-$  corresponds to the sign character of  $\mathcal{S}_{\psi_v}(G)$ . The following verifies the multiplicity formula for the global A-packet  $\Pi_{\psi}(G) = \bigoplus_v \Pi_{\psi_v}(G)$  conjectured by Arthur.

**Theorem 3.1.** If we write  $m(\pi)$  for the multiplicity of an irreducible representation  $\pi = \bigoplus_{v} \pi_{v} \in \prod_{\psi}(G)$  of  $G_{4}(\mathbb{A})$  in the L<sup>2</sup>-automorphic spectrum, then

$$m(\pi) = \frac{1}{|\mathcal{S}_{\psi}(G)|} \sum_{\bar{s} \in \mathcal{S}_{\psi}(G)} \epsilon_{\psi}(\bar{s}) \prod_{v} \langle \bar{s}, \pi_{v} \rangle.$$

Here

 $\epsilon_{\psi}(s) = \begin{cases} \operatorname{sgn}_{\mathcal{S}_{\psi}(G)} & \text{in the case } (2.b) \text{ with } \epsilon(1/2, \pi \times \eta) = -1, \\ 1 & \text{otherwise.} \end{cases}$ 

We show the inequality  $\geq$  by the global theta correspondence. Also in some cases, this gives the exact equality. But in other cases including (2.c), we need the analysis of the Fourier coefficients to have the converse inequality.

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