

The weak limit of renormalized Rabi Hamiltonian

Fumio Hiroshima

Faculty of Mathematics, Kyushu University, Japan

The 10th International Council for Industrial
and
Applied Mathematics

Aug.21,2023,Tokyo

Harmonic oscillators

$$a = \frac{1}{\sqrt{2}} \left(\frac{1}{\sqrt{\omega}} \frac{d}{dx} + \sqrt{\omega} x \right), \quad a^\dagger = \frac{1}{\sqrt{2}} \left(-\frac{1}{\sqrt{\omega}} \frac{d}{dx} + \sqrt{\omega} x \right)$$

are the annihilation op. and the creation op. in $L^2(\mathbb{R})$. They satisfy CCR

$$[a, a^\dagger] = \mathbb{1}.$$

Let

$$\omega a^\dagger a = \frac{1}{2} \left(-\frac{d^2}{dx^2} + \omega^2 x^2 - \omega \right).$$

Then $\varphi_g(x) = \left(\frac{\omega}{\pi}\right)^{1/4} e^{-\omega x^2/2}$ is the eigenvector associated with the lowest eigenvalue 0 of $\omega a^\dagger a$. The n th excited state of $\omega a^\dagger a$ is defined by

$$\varphi_n = \frac{1}{\sqrt{n!}} \prod_{k=0}^{n-1} a^\dagger \varphi_g, \quad n = 0, 1, 2, \dots$$

with $\varphi_0 = \varphi_g$. It follows that $\omega a^\dagger a \varphi_n = \omega n \varphi_n$ and $\text{Spec}(\omega a^\dagger a) = \{\omega n\}_{n=0}^\infty$.

$$L^2(\mathbb{R}) = \bigoplus_{n=0}^\infty [\mathbb{C} \varphi_n].$$

Spin 1/2

Let $\sigma = (\sigma_x, \sigma_y, \sigma_z)$ be elements of $SU(2)$.

$$\sigma_x = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}, \quad \sigma_y = \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix}, \quad \sigma_z = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$$

Let

$$\Delta\sigma_z, \quad \Delta > 0.$$

Then $\text{Spec}(\Delta\sigma_z) = \{\pm\Delta\}$.

The ground state of $\Delta\sigma_z$ is $s_- = \begin{bmatrix} 0 \\ 1 \end{bmatrix}$ and the excited state is

$s_+ = \begin{bmatrix} 1 \\ 0 \end{bmatrix}$. Then

$$\mathbb{C}^2 = \mathbb{C}s_+ \oplus \mathbb{C}s_-.$$

Rabi model

The total Hilbert space is $\mathbb{C}^2 \otimes L^2(\mathbb{R})$ and the free Rabi Hamiltonian is given by

$$\Delta\sigma_z \otimes \mathbb{1} + \mathbb{1} \otimes \omega a^\dagger a$$

Introducing the interaction $\sigma_x \otimes (a^\dagger + a)$, the Rabi Hamiltonian H_{Rabi} is given by

$$H_{\text{Rabi}} = \Delta\sigma_z \otimes \mathbb{1} + \mathbb{1} \otimes \omega a^\dagger a + g\sigma_x \otimes (a^\dagger + a), \quad g \in \mathbb{R}.$$

It has the parity symmetry:

$$[\sigma_z \otimes (-1)^{N/\omega}, H_{\text{Rabi}}] = 0.$$

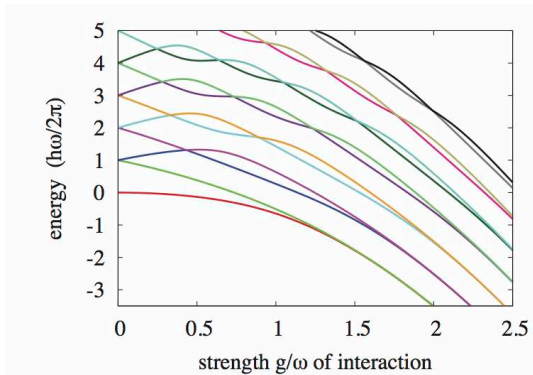


Figure: Eigenvalue curves

- 1) The lowest curve is concave.
- 2) Asymptotic behavior

Rotation U of spin

The rotation group in \mathbb{R}^3 has an adjoint representation on $SU(2)$. Let $n \in \mathbb{R}^3$ be a unit vector and $\theta \in [0, 2\pi)$. Thus $e^{(i/2)\theta n \cdot \sigma}$ satisfies that

$$e^{(i/2)\theta n \cdot \sigma} \sigma_\mu e^{-(i/2)\theta n \cdot \sigma} = (R\sigma)_\mu, \quad \mu = x, y, z$$

where $R = R(n, \theta)$ denotes 3×3 matrix representing the rotation around n with angle θ . In particular $U = e^{(i\pi/4)\sigma_y}$ yields that

$$U \sigma_x U^{-1} = \sigma_z,$$

$$U \sigma_z U^{-1} = -\sigma_x.$$

Then

$$U H_{\text{Rabi}} U^{-1} = -\Delta \sigma_x + \omega a^\dagger a + g \sigma_z (a^\dagger + a).$$

Ground state transformation U_g

Since φ_g is strictly positive, we can define the unitary operator

$$U_g : L^2(\mathbb{R}) \rightarrow L^2(\mathbb{R}, \varphi_g^2 dx), \quad U_g f = \varphi_g^{-1} f.$$

We set the probability measure on \mathbb{R} by

$$\varphi_g^2(x) dx = d\mu(x).$$

Thus

$$\begin{aligned} U_g U H_{\text{Rabi}} U^{-1} U_g^{-1} &= -\Delta \sigma_x - \frac{1}{2} \frac{d^2}{dx^2} + \omega x \frac{d}{dx} + g \sigma_z \sqrt{2\omega x} \\ &= \begin{pmatrix} -\frac{1}{2} \frac{d^2}{dx^2} + \omega x \frac{d}{dx} + g \sqrt{2\omega x} & -\Delta \\ -\Delta & -\frac{1}{2} \frac{d^2}{dx^2} + \omega x \frac{d}{dx} - g \sqrt{2\omega x} \end{pmatrix} \end{aligned}$$

Scalar transform

Let $\mathbb{Z}_2 = \{-1, 1\}$. We identify $\mathbb{C}^2 \otimes L^2(\mathbb{R}, d\mu)$ with

$$\mathcal{H} = L^2(\mathbb{R} \times \mathbb{Z}_2, d\mu) = \left\{ f = f(x, \alpha) \mid \sum_{\alpha \in \mathbb{Z}_2} \int |f(x, \alpha)|^2 d\mu(x) < \infty \right\}$$

by the map $\mathbb{C}^2 \otimes L^2(\mathbb{R}, d\mu) \ni \begin{bmatrix} f_+(x) \\ f_-(x) \end{bmatrix} \mapsto f(x, \alpha) \in \mathcal{H}$. Here $f(x, +1) = f_+(x)$ and $f(x, -1) = f_-(x)$. $H : \mathcal{H} \rightarrow \mathcal{H}$ is defined by

$$Hf(x, \alpha) = \left\{ -\frac{1}{2} \frac{d^2}{dx^2} + \omega x \frac{d}{dx} + g\sqrt{2\omega\alpha}x \right\} f(x, \alpha) - \Delta f(x, -\alpha), \quad \alpha \in \mathbb{Z}_2$$

Lemma

H_{Rabi} in $\mathbb{C}^2 \otimes L^2(\mathbb{R})$ is unitarily equivalent to H in $L^2(\mathbb{R} \times \mathbb{Z}_2, d\mu)$.

In what follows we deal with H instead of H_{Rabi} .

OU process

$(X_t^x)_{t \in \mathbb{R}}$ is an Ornstein-Uhlenbeck process on $(\mathcal{X}, \mathcal{B}, P)$ st

$$\mathbb{E}_P[X_t^x] = xe^{-\omega t}, \quad \mathbb{E}_P[X_t^x X_s^x] = e^{-\omega(t+s)} \left(x^2 + \frac{e^{2\omega(t \wedge s)} - 1}{2\omega} \right)$$

Here $\mathbb{E}_Q[\dots]$ denotes the expectation wrt Q . We have

$$\int \mathbb{E}_P[X_t^x] d\mu(x) = 0, \quad \int \mathbb{E}_P[X_t^x X_s^x] d\mu(x) = \frac{e^{-|t-s|\omega}}{2\omega}.$$

Let N be redefined by

$$N = -\frac{1}{2} \frac{d^2}{dx^2} + \omega x \frac{d}{dx} \quad (\text{the number op. in mind})$$

The generator of X_t^x is given by $-N$:

$$(f, e^{-tN} g)_{L^2(\mathbb{R}, d\mu)} = \int \mathbb{E}_P \left[\overline{f(X_0^x)} g(X_t^x) \right] d\mu(x).$$

Spin process

Let $(N_t)_{t \in \mathbb{R}}$ be a Poisson process on $(\mathcal{X}', \mathcal{B}', \nu)$ with unit intensity, i.e.,

$$\mathbb{E}_\nu [\mathbb{1}_{N_t=n}] = \frac{t^n}{n!} e^{-t}, \quad n \geq 0.$$

We define the spin process by

$$\sigma_t^\alpha = (-1)^{N_t} \alpha, \quad \alpha \in \mathbb{Z}_2.$$

Let σ_F be the fermionic harmonic oscillator:

$$\sigma_F = \frac{1}{2}(\sigma_z + i\sigma_y)(\sigma_z - i\sigma_y) = -\sigma_1 + \mathbb{1}.$$

The generator of σ_t^α is given by $-\sigma_F$:

$$(f, e^{-t\sigma_F} g)_{L^2(\mathbb{Z}_2)} = \sum_{\alpha \in \mathbb{Z}_2} \mathbb{E}_\nu \left[\overline{f(\sigma_0^\alpha)} g(\sigma_t^\alpha) \right].$$

Feynman-Kac Formula

Theorem (FH+Hirokawa, (14))

Set $\Sigma_{\alpha \in \mathbb{Z}_2} \int_{\mathbb{R}} \mathbb{E}_P \mathbb{E}_V [\dots] d\mu(x) = \mathbb{E}[\dots]$. Then

$$(f, e^{-tH} g)_{\mathcal{H}} = e^t \mathbb{E} \left[\overline{f(X_0^x, \sigma_0^\alpha)} g(X_t^x, \sigma_t^\alpha) e^{-g\sqrt{2\omega} \int_0^t \sigma_s^\alpha X_s^x ds} \Delta^{N_t} \right].$$

Corollary (FH+Hirokawa (14))

e^{-tH} is positivity improving for $\forall t > 0$. In particular $e^{-tH} \mathbb{1} > 0$.

Vacuum expectation

Let $\mathbb{1} \in L^2(\mathbb{R} \times \mathbb{Z}_2, d\mu)$ be the constant function. Note that

$$\mathbb{1} \cong \begin{bmatrix} 0 \\ 1 \end{bmatrix} \otimes \varphi_g \in \mathbb{C}^2 \otimes L^2(\mathbb{R}).$$

We set

$$\mathbb{1}_t = e^{-tH} \mathbb{1}, \quad t \geq 0.$$

Let $\Phi_g (> 0)$ be the ground state of H . Since

$$\Phi_g = \lim_{t \rightarrow \infty} \mathbb{1}_t / \|\mathbb{1}_t\|,$$

we have

$$(\Phi_g, A \Phi_g) = \lim_{t \rightarrow \infty} \frac{(\mathbb{1}_t, A \mathbb{1}_t)}{(\mathbb{1}_t, \mathbb{1}_t)}$$

Here $A = e^{+\beta N}, (-1)^N, \sigma_x(-1)^N$ etc.

$(B_t)_{t \geq 0}$ is Brownian motion. $\mathbb{E}[B_s B_t] = s \wedge t$.

$$X_s^x = e^{-s\omega} \left(x + \frac{1}{\sqrt{2\omega}} B_{e^{2s\omega}-1} \right) \quad s \geq 0$$

$$\begin{aligned} (\mathbb{1}, \mathbb{1}_t) &= e^t \sum_{\alpha \in \mathbb{Z}_2} \int \mathbb{E}_V \mathbb{E}_P \left[e^{-g\sqrt{2\omega} \int_0^t \sigma_s X_s^\alpha ds} \Delta^{N_t} \right] d\mu(x) \\ &= e^t \sum_{\alpha \in \mathbb{Z}_2} \mathbb{E}_V \left[\Delta^{N_t} \mathbb{E}_P \left[e^{-g \int_0^t \sigma_s^\alpha e^{-s\omega} B_{e^{2s\omega}-1} ds} \right] \int \sqrt{\frac{\omega}{\pi}} e^{-\omega x^2 - \sqrt{2\omega} g \int_0^t \sigma_s^\alpha e^{-s\omega} x dx} \right] \end{aligned}$$

$$\begin{aligned} 1) \mathbb{E}_P \left[e^{-g \int_0^t \sigma_s^\alpha e^{-s\omega} B_{e^{2s\omega}-1} ds} \right] &= \exp \left(\frac{g^2}{2} \left\| \int_0^t \sigma_s^\alpha e^{-s\omega} \mathbb{1}_{[0, e^{2s\omega}-1]}(\cdot) ds \right\|_{L^2(\mathbb{R})}^2 \right) \\ &= \exp \left(\frac{g^2}{2} \int_0^t ds \int_0^t dr \sigma_s^\alpha \sigma_r^\alpha e^{-(s+r)\omega} (e^{2(s \wedge r)\omega} - 1) \right), \end{aligned}$$

$$2) \sqrt{\frac{\omega}{\pi}} \int e^{-\omega x^2 - \sqrt{2\omega} g \int_0^t \sigma_s^\alpha e^{-s\omega} x dx} dx = \exp \left(\frac{g^2}{2} \left(\int_0^t \sigma_s^\alpha e^{-s\omega} ds \right)^2 \right)$$

Concave and parity

$$(\mathbb{1}, \mathbb{1}_t) = e^t \sum_{\alpha \in \mathbb{Z}_2} \mathbb{E}_v \left[\Delta^{N_t} \exp \left(\frac{g^2}{2} \int_0^t ds \int_0^t dr \sigma_s^\alpha \sigma_r^\alpha e^{-\omega|s-r|} \right) \right]$$

Let $E_0(g) = \inf \text{Spec}(H_{\text{Rabi}})$.

Theorem (Concave)

$g \mapsto E_0(g)$ is concave and differentiable a.e. in g .

Proof: This follows from $E_0(g) = \lim_{t \rightarrow \infty} -\frac{1}{t} \log(\mathbb{1}, \mathbb{1}_t)$. □

Theorem (Parity)

Let $P = \sigma_x \otimes (-1)^N$. Then $P\Phi_g = -\Phi_g$.

Proof: Since $P\mathbb{1} = -\mathbb{1}$ and $\Phi_g = \lim_{t \rightarrow \infty} \mathbb{1}_t / \|\mathbb{1}_t\|$, the theorem follows. □

$(\mathbb{1}_t, \mathbb{1}_t)$ and $(\mathbb{1}_t, e^{-\beta N} \mathbb{1}_t)$

$$\begin{aligned}(\mathbb{1}_t, \mathbb{1}_t) &= e^{2t} \sum_{\alpha \in \mathbb{Z}_2} \mathbb{E}_v \left[\Delta^{N_{2t}} \exp \left(\frac{g^2}{2} \int_0^{2t} ds \int_0^{2t} dr \sigma_s^\alpha \sigma_r^\alpha e^{-\omega|s-r|} \right) \right] \\ &= e^{2t} \sum_{\alpha \in \mathbb{Z}_2} \mathbb{E}_v \left[\Delta^{N_t + N_{-t}} \exp \left(\frac{g^2}{2} \int_{-t}^t ds \int_{-t}^t dr \sigma_s^\alpha \sigma_r^\alpha e^{-\omega|s-r|} \right) \right]\end{aligned}$$

Here we can extend $(N_t)_{t \geq 0}$ to $(N_t)_{t \in \mathbb{R}}$, where N_t is independent of N_{-s} for any $s, t > 0$, and we used the shift invariance of σ_t .

Since

$$(f, e^{-\beta N} g) = \sum_{\alpha \in \mathbb{Z}_2} \int \mathbb{E}_P[\bar{f}(X_0^x, \alpha) g(X_\beta^x, \alpha)] d\mu(x),$$

we see that

$$(\mathbb{1}_t, e^{-\beta N} \mathbb{1}_t) = \sum_{\alpha \in \mathbb{Z}_2} \int \mathbb{E}_P[\mathbb{1}_t(X_0^x, \alpha) \mathbb{1}_t(X_\beta^x, \alpha)] d\mu(x).$$

$$\begin{aligned} \mathbb{1}_t(X_0^x, \alpha) &= e^t \mathbb{E}_V \mathbb{E}_P \left[\Delta^{N_t} e^{-\sqrt{2\omega}g \int_0^t \sigma_s^\alpha X_s^x ds} \right] \\ &= e^t \mathbb{E}_V \mathbb{E}_P \left[\Delta^{N_t} e^{-\sqrt{2\omega}g \int_0^t \sigma_s^\alpha e^{-s\omega} \left(x + \frac{1}{\sqrt{2\omega}} B_{e^{2s\omega-1}} \right) ds} \right] \\ &= e^t \mathbb{E}_V \left[\Delta^{N_t} e^{-\sqrt{2\omega}g \int_0^t \sigma_s^\alpha e^{-s\omega} ds} \mathbb{E}_P \left[e^{-g \int_0^t \sigma_s^\alpha e^{-s\omega} B_{e^{2s\omega-1}} ds} \right] \right] \\ &= e^t \mathbb{E}_V \left[\Delta^{N_t} e^{-\sqrt{2\omega}g \int_0^t \sigma_s^\alpha e^{-s\omega} ds} e^{\frac{g^2}{2} \int_0^t ds \int_0^t dr \sigma_s^\alpha \sigma_r^\alpha e^{-(s+r)\omega} (e^{2(s\wedge r)\omega} - 1)} \right] \end{aligned}$$

$$\begin{aligned}
& \mathbb{E}_P \left[\mathbb{1}_t (X_\beta^x, \alpha) \right] \\
&= e^t \mathbb{E}_P \left[\mathbb{E}_V \left[\Delta^{N_t} e^{-\sqrt{2}\omega g \int_0^t \sigma_s^\alpha e^{-s\omega} ds} X_\beta^x e^{\frac{g^2}{2} \int_0^t ds \int_0^t dr \sigma_s^\alpha \sigma_r^\alpha e^{-(s+r)\omega} (e^{2(s\wedge r)\omega} - 1)} \right] \right] \\
&= e^t \mathbb{E}_P \left[\mathbb{E}_V \left[\Delta^{N_t} e^{-\sqrt{2}\omega g \int_0^t \sigma_s^\alpha e^{-s\omega} ds} e^{-\beta\omega} \left(x + \frac{1}{\sqrt{2}\omega} B_{e^{2\beta\omega} - 1} \right) \right. \right. \\
&\quad \left. \left. \times e^{\frac{g^2}{2} \int_0^t ds \int_0^t dr \sigma_s^\alpha \sigma_r^\alpha e^{-(s+r)\omega} (e^{2(s\wedge r)\omega} - 1)} \right] \right] \\
&= e^t \mathbb{E}_V \left[\Delta^{N_t} e^{-\sqrt{2}\omega g \int_0^t \sigma_s^\alpha e^{-s\omega} ds} e^{-\beta\omega} x \right. \\
&\quad \left. \times e^{\frac{g^2}{2} \int_0^t ds \int_0^t dr \sigma_s^\alpha \sigma_r^\alpha e^{-(s+r)\omega} (1 - e^{-2\beta\omega})} e^{\frac{g^2}{2} \int_0^t ds \int_0^t dr \sigma_s^\alpha \sigma_r^\alpha e^{-(s+r)\omega} (e^{2(s\wedge r)\omega} - 1)} \right]
\end{aligned}$$

Let $W(s, r) = \sigma_s^\alpha \sigma_r^\alpha e^{-\omega|s-r|}$. Finally we have

$$\begin{aligned} & (\mathbb{1}_t, e^{-\beta N} \mathbb{1}_t) \\ &= e^{2t} \sum_{\alpha \in \mathbb{Z}_2} \mathbb{E}_v \left[\Delta^{N_{-t} + N_t} \exp \left(\frac{g^2}{2} \int_{-t}^t \int_{-t}^t W(s, r) ds dr \right) \right. \\ & \quad \left. \times \exp \left(-g^2 (1 - e^{-\beta \omega}) \int_{-t}^0 \int_0^t W(s, r) ds dr \right) \right] \end{aligned}$$

Note that

$$\begin{aligned} \sigma_s^\alpha \sigma_r^\alpha &= (-1)^{N_s} (-1)^{N_r} \alpha^2 \\ &= (-1)^{N_s} (-1)^{N_r} = (-1)^{N_s} (-1)^{-N_r} = (-1)^{N_s - N_r} = (-1)^{N_{|s-r|}}, \end{aligned}$$

which implies that $\sigma_s^\alpha \sigma_r^\alpha$ is independent of $\alpha \in \mathbb{Z}_2$. Thus

$$e^{i\pi N_{|s-r|}} = \sigma_s^\alpha \sigma_r^\alpha$$

Let $W(s-r) = e^{i\pi N_{|s-r|}} e^{-\omega|s-r|}$.

$$\begin{aligned}
 (\Phi_g, e^{-\beta N} \Phi_g) &= \lim_{t \rightarrow \infty} \frac{(\mathbb{1}_t, e^{-\beta N} \mathbb{1}_t)}{(\mathbb{1}_t, \mathbb{1}_t)} \\
 &= \lim_{t \rightarrow \infty} \frac{\sum_{\alpha \in \mathbb{Z}_2} \mathbb{E}_{\mathbf{V}} \left[\Delta^{N_t + N_{-t}} e^{\frac{g^2}{2} \int_{-t}^t \int_{-t}^t W(s-r) ds dr} - g^2 (1 - e^{-\beta \omega}) \int_{-t}^0 \int_0^t W(s-r) ds dr \right]}{\sum_{\alpha \in \mathbb{Z}_2} \mathbb{E}_{\mathbf{V}} \left[\Delta^{N_t + N_{-t}} e^{\frac{g^2}{2} \int_{-t}^t \int_{-t}^t W(s-r) ds dr} \right]} \\
 &= \lim_{t \rightarrow \infty} \mathbb{E}_{\mathbf{V}_t} \left[e^{-g^2 (1 - e^{-\beta \omega}) \int_{-t}^0 \int_0^t W(s-r) ds dr} \right] \\
 \nu_t(\mathcal{O}) &= \sum_{\alpha \in \mathbb{Z}_2} \mathbb{E}_{\mathbf{V}} \left[\Delta^{N_t + N_{-t}} e^{\frac{g^2}{2} \int_{-t}^t \int_{-t}^t W(s-r) ds dr} \mathcal{O} \right] / \sum_{\alpha \in \mathbb{Z}_2} \mathbb{E}_{\mathbf{V}} \left[\Delta^{N_t + N_{-t}} e^{\frac{g^2}{2} \int_{-t}^t \int_{-t}^t W(s-r) ds dr} \right]
 \end{aligned}$$

Theorem (FH Adv. Math. (14))

Let $W(s-r) = e^{i\pi N_{|s-r|}} e^{-\omega|s-r|}$. There exists a prob. measure ν_∞ st

$$(\Phi_g, e^{-\beta N} \Phi_g) = \mathbb{E}_{\nu_\infty} \left[e^{-g^2 (1 - e^{-\beta \omega}) \int_{-\infty}^0 \int_0^\infty W(s-r) ds dr} \right] \quad \beta > 0$$

Localization

$\mathbb{C}^2 \otimes L^2(\mathbb{R}) = \bigoplus_{n=0}^{\infty} [\mathbb{C}^2 \otimes [\mathbb{C}\varphi_n]]$. Thus $\Phi_g = \bigoplus_{n=0}^{\infty} \Phi_g^{(n)}$ and

$$1 = \|\Phi_g\|^2 = \sum_{n=0}^{\infty} \|\Phi_g^{(n)}\|_{\mathbb{C}^2 \otimes [\mathbb{C}\varphi_n]}^2$$

The question is $\|e^{\beta N} \Phi_g\|^2 = \sum_{n=0}^{\infty} e^{2\beta n} \|\Phi_g^{(n)}\|_{\mathbb{C}^2 \otimes [\mathbb{C}\varphi_n]}^2 < \infty?$

Corollary (Localization)

Let $W(s-r) = e^{i\pi N_{|s-r|}} e^{-\omega|s-r|}$. By the analytic continuation on β we see that

$$(\Phi_g, e^{\beta N} \Phi_g) = \mathbb{E}_{V_\infty} \left[e^{-g^2(1-e^{\beta\omega}) \int_{-\infty}^0 \int_0^\infty W(s-r) ds dr} \right] \quad \beta \in \mathbb{C},$$

$$(\Phi_g, N \Phi_g) = \omega g^2 \mathbb{E}_{V_\infty} \left[\int_{-\infty}^0 \int_0^\infty W(s-r) ds dr \right]$$

In particular

$$\|e^{\beta N} \Phi_g\|^2 < \infty, \quad \forall \beta > 0.$$

$\sum_{n=0}^{\infty} \|\Phi_g^{(n)}\|^2 = 1$, but

$$(\Phi_g, (-1)^{N/\omega} \Phi_g) = \sum_{n=0}^{\infty} (-1)^n \|\Phi_g^{(n)}\|^2 \leq 0?$$

Theorem

$$(\Phi_g, (-1)^{N/\omega} \Phi_g) > 0$$

Proof: Let $\beta = i\pi/\omega$. Then

$$(\Phi_g, (-1)^{N/\omega} \Phi_g) = \mathbb{E}_{V_\infty} \left[e^{2g^2 \int_{-\infty}^0 \int_0^\infty e^{i\pi N|s-r|} e^{-\omega|s-r|} ds dr} \right] > 0$$



Spectral zetafunction

Let $\omega = 1$ and $p = -id/dx$.

$$H_{\text{Rabi}} + g^2 \cong H_{\text{ren}} = \begin{pmatrix} a^\dagger a & 0 \\ 0 & a^\dagger a \end{pmatrix} - \Delta \begin{pmatrix} 0 & e^{-i2\sqrt{2}gp} \\ e^{+i2\sqrt{2}gp} & 0 \end{pmatrix}$$

Let $\text{Spec}(H_{\text{Rabi}}) = \{E_n(g)\}_{n=0}^\infty$.

If $\Delta < 1/4$, then $E_{2n}, E_{2n+1} \in (n - 1/4, n + 1/4)$. Let

$$\zeta_g(s) = \sum_{n=0}^{\infty} \frac{1}{(E_n(g) + g^2)^s}.$$

$$\text{Let } H = \begin{pmatrix} a^\dagger a & 0 \\ 0 & a^\dagger a \end{pmatrix}.$$

Lemma

Let $\tau > 0$ and $s > 0$. Then $s - \lim_{g \rightarrow \infty} (H_{\text{ren}} + \tau)^{-s} = (H + \tau)^{-s}$.

Theorem

$$\lim_{g \rightarrow \infty} \sum_{n=0}^{\infty} \frac{1}{(E_n(g) + g^2 + \tau)^s} = 2 \sum_{n=0}^{\infty} \frac{1}{(n + \tau)^s}$$

Proof: Let $e_{2n} = \begin{bmatrix} \varphi_n \\ 0 \end{bmatrix}$ and $e_{2n+1} = \begin{bmatrix} 0 \\ \varphi_n \end{bmatrix}$.

$$\begin{aligned} \lim_{g \rightarrow \infty} \zeta_g(s) &= \lim_{g \rightarrow \infty} \text{Tr}(H_{\text{ren}} + \tau)^{-s} = \lim_{g \rightarrow \infty} \sum_{n=0}^{\infty} (e_n, (H_{\text{ren}} + \tau)^{-s} e_n) \\ &= \sum_{n=0}^{\infty} (e_n, (H + \tau)^{-s} e_n) \end{aligned}$$

□

The same result is obtained in

Cid Reyes-Bustos and Masato Wakayama, arXiv:2304.08943, 2023.

Summary

- $(f, e^{-tH}g) = e^t \sum_{\alpha \in \mathbb{Z}_2} \int_{\mathbb{R}} \mathbb{E}_P \mathbb{E}_V \left[\overline{f(X_0^x, \sigma_0^\alpha)} g(X_t^x, \sigma_t^\alpha) e^{-g\sqrt{2\omega} \int_0^t \sigma_s^\alpha X_s^x ds} \Delta^{N_t} \right] d\mu(x).$
- $(\mathbb{1}, e^{-tH} \mathbb{1}) = e^t \sum_{\alpha \in \mathbb{Z}_2} \mathbb{E}_V \left[\Delta^{N_t} \exp \left(\frac{g^2}{2} \int_0^t ds \int_0^t dr e^{i\pi N_{|s-r|}} e^{-\omega|s-r|} \right) \right]$
- Prob. meas. $\exists v_\infty$ st $(\Phi_g, e^{\beta N} \Phi_g) = \mathbb{E}_{v_\infty} \left[e^{-g^2(1-e^{\beta\omega}) \int_{-\infty}^0 \int_0^\infty e^{i\pi N_{|s-r|}} e^{-\omega|s-r|} ds dr} \right] \quad \forall \beta \in \mathbb{C}.$
- $g \mapsto E_0(g)$ is concave and differentiable a.e. in g .
- $\sigma_x(-1)^{N/\omega} \Phi_g = -\Phi_g.$
- $\|e^{\beta N} \Phi_g\|^2 < \infty, \quad \forall \beta > 0.$
- $\lim_{g \rightarrow \infty} \sum_{n=0}^{\infty} \frac{1}{(E_n(g) + g^2 + \tau)^s} = 2 \sum_{n=0}^{\infty} \frac{1}{(n + \tau)^s}$

Thank you.