# Sharp interface limit of stochastic Cahn-Hilliard equation

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### Joint work with Lubomir Banas and Huanyu Yang

[Banas/Yang/Z.: arXiv:1905.09182] [Yang/Z.:arXiv:1905.07216]



2 Singular noise for large  $\sigma$ 



(3) Weak approach and tightness for small  $\sigma$ 

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The minimizers of the energy (2) are the constant functions  $u \equiv 1$  and  $u \equiv -1$ , which represent the "pure phases" of the system.

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• As  $\varepsilon \to 0$ ,  $F(u^{\varepsilon}) \to 0$ , which implies  $u^{\varepsilon} \to -1 + 2\mathbf{1}_E$  for some  $E \subset [0, T] \times \mathcal{D}$ .  $\Gamma_t := \partial E_t$  is the interface.

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- Formally,

$$\begin{split} \int_0^t \int_{\mathcal{D}} \partial_t \mathbf{1}_{E_t} \psi &= -\frac{1}{2} \int_0^t \int_{\mathcal{D}^+} \nabla v \nabla \psi - \frac{1}{2} \int_0^t \int_{\mathcal{D}^-} \nabla v \nabla \psi \\ &= \frac{1}{2} \int_0^t \int_{\mathcal{D}^+} \operatorname{div}(\nabla v \psi) + \frac{1}{2} \int_0^t \int_{\mathcal{D}^-} \operatorname{div}(\nabla v \psi) \\ &= \frac{1}{2} \int_0^t \int_{\Gamma_t} (\partial_n v^+ - \partial_n v^-) \psi. \end{split}$$

# Hele-Shaw model

Formally derived by [Pego: 1989] and rigorous proved by [Alikakos, Bates, Chen: 1994]:  $v^{\varepsilon} \rightarrow v$ ,  $(v, \Gamma)$  solves the following free boundary problem:

$$\begin{split} \Delta v &= 0 \text{ in } \mathcal{D} \setminus \Gamma_t, \ t > 0, \\ \frac{\partial v}{\partial n} &= 0 \text{ on } \partial \mathcal{D}, \\ v &= \frac{2}{3}H \text{ on } \Gamma_t, \\ \mathcal{V} &= \frac{1}{2}(\partial_n v^+ - \partial_n v^-) \text{ on } \Gamma_t, \end{split}$$

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**Problems**: Singular noise or Small  $\sigma \Rightarrow$ ? Stochastic Hele shaw

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• Stopping time argument yields that

#### Theorem 1

[Banas, Yang, Z. 19] For  $\sigma > \frac{107}{12}$ ,  $\|R^{\varepsilon}\|_{L^{3}_{t}L^{3}_{x}}$  converges to 0 in probability.

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• Da prato-Debussche's trick:  $Z_t^{\varepsilon} := \varepsilon^{\sigma} \int_0^t e^{-(t-s)\varepsilon\Delta^2} dW_s$ ,  $Y^{\varepsilon} := R^{\varepsilon} - Z^{\varepsilon}$  satisfies:

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- $\mathbb{E}(\|Z^{\varepsilon}\|_{\mathcal{C}(\mathcal{D}_{T})}) \lesssim \varepsilon^{(\sigma \frac{1}{4})^{-}};$
- Stopping time argument yields that

#### Theorem 1

[Banas, Yang, Z. 19] For  $\sigma > \frac{107}{12}$ ,  $||R^{\varepsilon}||_{L_t^3 L_x^3}$  converges to 0 in probability.  $\Rightarrow$  the limit is the same as the deterministic case

$$du^{\varepsilon,h} = \Delta\left(-\varepsilon\Delta u^{\varepsilon,h} + \frac{1}{\varepsilon}\left(F'(u^{\varepsilon,h}) - 3\frac{c_{h,t}^{\varepsilon}}{c_{h,t}^{\varepsilon}}u^{\varepsilon,h}\right)\right)dt + \varepsilon^{\sigma}\nabla\cdot dW_{t}^{h},$$

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#### Theorem 2

[Banas, Yang, Z. 19] Assume  $\varepsilon^{\theta} \lesssim h^2$ . Then for  $\sigma > \frac{26}{3} + \theta$ ,  $\|R^{\varepsilon,h}\|_{L^3_t L^3_x}$  converges to 0 in probability.

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We consider the stochastic Cahn-Hilliard equation on a bounded smooth open domain  $\mathcal{D} \subset \mathbb{R}^d$  (d = 2, 3):

$$\begin{cases}
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# Idea of Proof: Lyapunov property

Recall

$$\mathcal{E}^{\varepsilon}(u^{\varepsilon}) := rac{\varepsilon}{2} \int_{\mathcal{D}} |\nabla u^{\varepsilon}(x)|^2 dx + rac{1}{\varepsilon} \int_{\mathcal{D}} F(u^{\varepsilon}(x)) dx.$$

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$$= - \langle \nabla v^{\varepsilon}, \nabla v^{\varepsilon} \rangle dt + \frac{\varepsilon^{2\sigma+1}}{2} \operatorname{Tr}(-\Delta Q)dt + \frac{\varepsilon^{2\sigma-1}}{2} \operatorname{Tr}(F''(u^{\varepsilon})Q)dt + \varepsilon^{\sigma} \langle v^{\varepsilon}, dW_{t} \rangle$$

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#### Lemma 3

(Lyapunov property) Assume  $Tr(-\Delta Q) < \infty$  and  $\sup_{0 < \varepsilon < 1} \mathcal{E}^{\varepsilon}(u_0^{\varepsilon}) < \mathcal{E}_0$  then there exists  $\varepsilon_0 \in (0, 1)$  such that for any  $\varepsilon \in (0, \varepsilon_0]$  and any  $p \ge 1$ ,

$$\mathbb{E} \sup_{t\in[0,T]} \mathcal{E}^{arepsilon}(t)^{
ho} \lesssim (arepsilon^{2\sigma-1}+\mathcal{E}_0)^{
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Rongchan Zhu (Beijing Institute of Technology)

Idea of Proof: Tightness for  $\sigma \geq \frac{1}{2}$ 

#### Lemma 4

# Assume $\sigma \geq \frac{1}{2}$ . For any $\beta \in (0, \frac{1}{12})$ $\mathbb{E} \left( \| u^{\varepsilon} \|_{C^{\beta}([0,T];L^{2})} \right) \lesssim 1$

For any  $\delta > 0$ , there exists a constant  $C \equiv C(\delta, T) > 0$ , such that

$$\mathbb{P}\left(\int_0^{\mathcal{T}} \| v^arepsilon(t) \|_{H^1}^2 dt \leq C
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 $\mathsf{Tightness} + \mathsf{Skorohod \ theorem} \Rightarrow$ 

#### Theorem 5

[Y, Zhu: 19] For any 
$$\sigma \ge \frac{1}{2}$$
,  $\mathbb{P} - a.s. \omega$ ,  
•  $u^{\varepsilon} \rightarrow -1 + 2\mathbf{1}_{E}$  in  $C([0, T], L^{2}_{w})$ ,  
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 $\Rightarrow$  rigorously proved in radial symmetric case and conjectured in general case to the deterministic Hele Shaw model:

$$\begin{cases} \Delta v = 0 \text{ in } \mathcal{D} \setminus \Gamma_t, \ t > 0, \\ \frac{\partial v}{\partial n} = 0 \text{ on } \partial \mathcal{D}, \\ v = \frac{2}{3}H \text{ on } \Gamma_t, \\ \mathcal{V} = \frac{1}{2}(\partial_n v^+ - \partial_n v^-) \text{ on } \Gamma_t \end{cases}$$

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 $\Rightarrow$  The only possibility to converge to "stochastic Hele-Shaw model" is  $\sigma = 0$ .

(9)

#### Weak approach and tightness for small $\sigma$

## Equation driven by "smeared" noise

We consider the following random PDE:

$$\begin{cases} \frac{\partial u^{\varepsilon}}{\partial t} = \Delta v^{\varepsilon} + \varepsilon^{\sigma} \xi^{\varepsilon}_{t}, \quad (t, x) \in [0, T] \times \mathcal{D}, \\ v^{\varepsilon} = -\varepsilon \Delta u^{\varepsilon}(t) + \frac{1}{\varepsilon} F'(u^{\varepsilon}(t)), \quad (t, x) \in [0, T] \times \mathcal{D}, \\ \frac{\partial u^{\varepsilon}}{\partial n} = \frac{\partial v^{\varepsilon}}{\partial n} = 0, \quad (t, x) \in [0, T] \times \partial \mathcal{D}, \\ u^{\varepsilon}(0, x) = u^{\varepsilon}_{0}(x), \quad x \in \mathcal{D}, \end{cases}$$
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 deterministic model (9).
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Main results for  $\sigma = 0$  [Y, Zhu: 19]

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• 
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• Let  $\hat{v} = v + \Delta^{-1}\xi$   

$$\int_0^t \int_{\mathcal{D}} \partial_t \mathbf{1}_{E_t} \psi = -\frac{1}{2} \int_0^t \int_{\mathcal{D}^+} \nabla \hat{v} \nabla \psi - \frac{1}{2} \int_0^t \int_{\mathcal{D}^-} \nabla \hat{v} \nabla \psi$$

$$= \frac{1}{2} \int_0^t \int_{\Gamma_t} (\partial_n \hat{v}^+ - \partial_n \hat{v}^-) \psi.$$

For  $\sigma = 0$ , we proved in radial symmetric case, that the sharp interface limit of (10) is the weak formula of the following "stochastic Hele-Shaw" model:

$$\begin{cases} \Delta v dt = -dW_t \text{ in } \mathcal{D} \setminus \Gamma_t, \ t > 0, \\ \frac{\partial v}{\partial n} = 0 \text{ on } \partial \mathcal{D}, \\ v = \frac{2}{3} H \text{ on } \Gamma_t, \\ \mathcal{V} dt = \frac{1}{2} \left[ \frac{\partial}{\partial n} \right]_{\Gamma_t} (v dt + \Delta^{-1} dW_t), \end{cases}$$

(11)

where H: mean curvature

$$\left[\frac{\partial}{\partial n}\right]_{\Gamma_t} f = \frac{\partial f^+}{\partial n} - \frac{\partial f^-}{\partial n}.$$

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N D Alikakos, P W Bates, and Xinfu Chen. Convergence of the Cahn-Hilliard Equation to the Hele-Shaw Model. *Archive for Rational Mechanics and Analysis*, 128(2):165–205, 1994.

## D C Antonopoulou, D Blömker, and G D Karali. The sharp interface limit for the stochastic Cahn-Hilliard equation. *Annales de l'Institut Henri Poincaré Probabilités et Statistiques*, 54(1):280–298, 2018.

Lubomir Banas, Huanyu Yang, and Rongchan Zhu. Sharp interface limit of stochastic Cahn-Hilliard equation with singular noise. *arXiv.org*, May 2019.



#### Xinfu Chen.

Global asymptotic limit of solutions of the Cahn-Hilliard equation. Journal of Differential Geometry, 44(2):262–311, 1996.

Giuseppe Da Prato and Arnaud Debussche. Stochastic Cahn-Hilliard equation.

Nonlinear Analysis: Theory, Methods & Applications, 26(2):241–263, 1996.

#### J E Hutchinson and Y Tonegawa.

Convergence of phase interfaces in the van der Waals-Cahn-Hilliard theory. *Calculus of Variations and Partial Differential Equations*, 10(1):49–84, January 2000.

R L Pego.

Front Migration in the Nonlinear Cahn-Hilliard Equation.

Proceedings of the Royal Society A: Mathematical, Physical and Engineering Sciences, 422(1863):261–278, April 1989.



Michael Röckner, Huanyu Yang, and Rongchan Zhu. Conservative stochastic 2-dimensional Cahn-Hilliard equation. February 2018.



Matthias Röger and Reiner Schaetzle. On a modified conjecture of De Giorgi. *Mathematische Zeitschrift*, 254(4):675–714, December 2006.

Matthias Röger and Yoshihiro Tonegawa. Convergence of phase-field approximations to the Gibbs-Thomson law. *Calculus of Variations and Partial Differential Equations*, 32(1):111–136, 2008.

# Thank you!