# Derivation of viscous Burgers equations from weakly asymmetric simple exclusion processes

Kenkichi TSUNODA joint with M. Jara and C. Landim

Department of Mathematics, Osaka University

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## Outline of the talk

- Overview and Related Works
- 2 Model and Main Results
- Comments on the Proof

## 1. Overview and Related Works

### Overview

■ The macroscopic density of the  $(WASEP)_n$  evolves according to the nonlinear heat eq. as the system size n grows to infinity (Hydrodynamic limit):

$$\partial_t u = \nabla \cdot [D(u)\nabla u] + \nabla \cdot [\sigma(u)\mathbf{m}] , \qquad (1)$$

where  $D, \sigma$  are  $d \times d$ -matrices (Diffusivity and Mobility).

■ For small  $\varepsilon > 0$ , let us consider the first order correction to (1) around a constant profile  $\alpha_0 \in (0,1)$ :

$$\begin{cases} \partial_t u^{\varepsilon} = \nabla \cdot [D(u^{\varepsilon}) \nabla u^{\varepsilon}] + \varepsilon^{-1} \nabla \cdot [\sigma(u^{\varepsilon}) \mathbf{m}] , \\ u^{\varepsilon}(0, \cdot) = \alpha_0 + \varepsilon v_0 , \end{cases}$$
 (2)

for some smooth function  $v_0$ .

- The solution  $u^{\varepsilon}$  should evolve as  $u_t^{\varepsilon} \sim \alpha_0 + \varepsilon v_t$ .
- Indeed, if  $\sigma'(\alpha_0) = 0$ , the sequence  $\{\varepsilon^{-1}(u^{\varepsilon} \alpha_0)\}_{\varepsilon>0}$  converges to the solution to the Burgers eq. as  $\varepsilon \downarrow 0$  (Incompressible limit):

$$\begin{cases} \partial_t v = \nabla \cdot [D(\alpha_0) \nabla v] + (1/2) \nabla \cdot [v^2 \sigma''(\alpha_0) \mathbf{m}] , \\ v(0, \cdot) = v_0(\cdot) . \end{cases}$$

■ Main Result (rough version): Taking  $\varepsilon = \varepsilon_n \downarrow 0 \ (n \to \infty)$ , the correctly scaled density of the WASEP evolves according to the Burgers eq.

## Related Works

- Many results on hydrodynamic limits.
   e.g. Guo-Papanicolaou-Varadhan, 88, Yau, 91.
- Esposito-Marra-Yau, 94, 96 · · · Derivation of Burgers equation and Navier-Stokes equation  $(d \ge 3)$ .
- Quastel-Yau,  $98 \cdots$  Large deviations for the incompressible limits (d = 3).
- Beltán-Landim, 08 · · · Derivation of Burgers equation and Navier-Stokes equation in any dimensions but with (meso-scopically) big jumps.
- Jara-Menezes, 19+ · · · Sharp entropy bound.

## 2. Model and Main Results

## Model

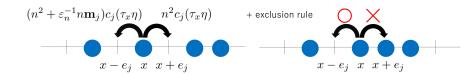
- Each particle moves on the *d*-dimensional discrete torus  $\mathbb{T}_n^d = (\mathbb{Z}/n\mathbb{Z})^d = \{1, 2, \cdots, n\}^d, n \in \mathbb{N}$ . Let  $\mathbb{T}^d$  be the *d*-dimensional torus  $\mathbb{T}^d = (\mathbb{R}/\mathbb{Z})^d = [0, 1)^d$ .
- Denote the number of particles at site  $x \in \mathbb{T}_n^d$  at time t by  $\eta_t^n(x)$   $(\eta_t^n = \{\eta_t^n : x \in \mathbb{T}_n^d\} \in \{0,1\}^{\mathbb{T}_n^d})$ .
- Some parameters:
  - $\{\varepsilon_n\}_{n\in\mathbb{N}}\subset\mathbb{R}$ : a sequence converging to 0.
  - $(c_j)_{j=1}^d$ : nonnegative local functions.
  - $\mathbf{m} = (\mathbf{m}_1, \dots, \mathbf{m}_d)$ : a vector in  $\mathbb{R}^d$ .

Let  $\eta_t^n$  be a Markov process on  $\{0,1\}^{\mathbb{T}_n^d}$  with the generator  $L_n=n^2L_n^S+\varepsilon_n^{-1}nL_n^A$  with

for a function  $f: \{0,1\}^{\mathbb{T}_n^d} \to \mathbb{R}$ .

The dynamics of our particle system is as follows:

■ Each particle can jump from x to x+1 or x-1 at given rates only if the site x is occupied and the site x+1 or x-1 is vacant.



■ For a continuous function  $u_0 : \mathbb{T}^d \to [0,1]$ , let  $\nu_{u_0}^n$  be the product Bernoulli measure on  $\{0,1\}^{\mathbb{T}_n^d}$ :

$$u_{u_0}^n (\eta : \eta(x) = 1) = u_0(x/n), \quad x \in \mathbb{T}_n^d.$$

■ Gradient condition: For each j, there exist finitely supported signed measures  $m_{j,p}$ ,  $p=1,\ldots,n_j$  and local functions  $g_{j,p}$  such that

$$c_{j}(\eta)[\eta_{0} - \eta_{e_{j}}] = \sum_{p=1}^{n_{j}} \sum_{y \in \mathbb{Z}^{d}} m_{j,p}(y) g_{j,p}(\tau_{y}\eta) ,$$
  
 $\sum_{y \in \mathbb{Z}^{d}} m_{j,p}(y) = 0 .$ 

# Classical case $(\varepsilon_n = 1)$

- Assume that  $\eta_0^n \stackrel{\mathrm{d}}{=} \nu_{u_0}^n$  for some continuous function  $u_0 : \mathbb{T}^d \to [0,1]$ .
- Hydrodynamic limit: For any  $t \ge 0$  and any smooth function  $H: \mathbb{T}^d \to \mathbb{R}$ ,

$$\lim_{n\to\infty} \mathbb{E}^n \left[ \left| \frac{1}{n^d} \sum_{x\in \mathbb{T}_n^d} H(x/n) \eta_t^n(x) - \int_{\mathbb{T}^d} H(x) u(t,x) dx \right| \right] = 0,$$

where  $u:[0,\infty)\times\mathbb{T}^d\to[0,1]$  is the unique weak solution of the Cauchy problem

$$\begin{cases} \partial_t u = \nabla \cdot [D(u)\nabla u] + \nabla \cdot [\sigma(u)\mathbf{m}] , \\ u(0,\cdot) = u_0(\cdot) . \end{cases}$$

# Incompressible case $(\varepsilon_n \downarrow 0)$

- Fix  $\alpha_0 \in (0,1)$  with  $\sigma'(\alpha_0) = 0$  and assume that  $\eta_0^n \stackrel{\mathrm{d}}{=} \nu_{\alpha_0 + \varepsilon_n \nu_0}^n$  for some function  $\nu_0 \in C^{3+}(\mathbb{T}^d)$ .
- Let  $v:[0,\infty)\times\mathbb{T}^d\to\mathbb{R}$  be the unique weak (classical) solution of the Burgers eq.:

$$\begin{cases} \partial_t v = \nabla \cdot [D(\alpha_0) \nabla v] + (1/2) \nabla \cdot [v^2 \sigma''(\alpha_0) \mathbf{m}] , \\ v(0, \cdot) = v_0(\cdot) . \end{cases}$$

■ For each  $t \ge 0$ , let  $u_t^n = \alpha_0 + \varepsilon_n v_t$ ,  $\nu_t^n = \nu_{u_t^n}^n$  and let  $\mu_t^n$  be the distribution of  $\eta_t^n$ .

## Main Results

#### Theorem 1 (Jara-Landim-T., 19+)

Assume that  $n^2 \varepsilon_n^4 \leq C_0 g_d(n)$  for some constant  $C_0$ , where  $g_d(n) = n, \log n, 1$  if  $d = 1, d = 2, d \geq 3$ , respectively. Then, for any T > 0, there exists a constant  $C_1 = C_1(T, v_0, C_0)$  such that for any  $0 \leq t \leq T$ ,

$$H(\mu_t^n|\nu_t^n) \leq C_1 n^{d-2} g_d(n) ,$$

where  $H(\mu|\nu)$  is the relative entropy of  $\mu$  w.r.t.  $\nu$ .

## Corollary 2 (Jara-Landim-T., 19+)

Assume that  $n^2 \varepsilon_n^4 \leq C_0 g_d(n)$  and  $\varepsilon_n^2 n^2 g_d(n)^{-1} \uparrow \infty$ . For any  $t \geq 0$  and any smooth function  $H : \mathbb{T}^d \to \mathbb{R}$ ,

$$\lim_{n\to\infty} \mathbb{E}^n \left[ \left| \frac{1}{\varepsilon_n n^d} \sum_{x\in \mathbb{T}_n^d} H(x/n) [\eta_t^n(x) - \alpha_0] - \int_{\mathbb{T}^d} H(x) v(t,x) dx \right| \right] = 0,$$

#### Remarks

- Initial distribution: The assumption  $\eta_0^n \stackrel{\mathrm{d}}{=} \nu_{\alpha_0 + \varepsilon_n \nu_0}^n$  can be replaced with the entropy bound at time 0.
- $\sigma'(\alpha_0) = 0$ : In the case of general  $\alpha \in (0,1)$ , introducing the Galilean transformation  $\alpha + \varepsilon_n v(t, x \varepsilon_n^{-1} \sigma'(\alpha) \mathbf{m} t)$ , we can obtain a similar result.

## 3. Comments on the Proof

■ Following [Jara-Menezes, 19+], we shall compute the entropy production. Let  $H_t = H(\mu_t^n | \nu_t^n)$ . Then, we have

$$\frac{d}{dt}H_t \leq -n^2D(g_t^n, L_n^S, \nu_t^n) + \int \left\{L_n^{*,\nu_t^n}\mathbf{1} - \partial_t \log \nu_t^n\right\} d\mu_t^n.$$

- We need to compute the integrand  $L_n^{*,\nu_t^n}\mathbf{1} \partial_t \log \nu_t^n$  explicitly. Indeed, it can be expressed in terms of the "Fourier coefficients" of  $g_{j,p}$  (but quite messy...).
- We also need to expand several terms in  $\varepsilon_n$  properly: e.g.  $E_{\nu_n^n}[\tau_{-e_i}g_{j,p}-g_{j,p}]$ .
- Together with these calculations, we shall apply techniques developed in [Jara-Menezes, 19+].

Thank you for your attention.